

(July 3, 2006)

## Poles of half-degenerate Eisenstein series

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- Review of a well-known example
- Simplest non-self-associate functional equation
- Half-cuspidal data
- Associate Eisenstein series

We study the meromorphic continuation and functional equations for a useful special class of half-degenerate Eisenstein series for maximal proper parabolics

$$P = P^{r-q,q} = \left\{ \begin{pmatrix} (r-q)\text{-by-}(r-q) & * \\ 0 & q\text{-by-}q \end{pmatrix} \right\}$$

on  $GL_r$  over a number field  $k$ , involving *cuspidal data* on the lower-right  $GL_q$  factor of the (standard) Levi component and *non-cuspidal data* on the upper-left  $GL_{r-q}$  factor. In particular, with a normalizing factor of a (single) Godement-Jacquet zeta integral attached to the cuspidal data, *there are no poles*.

As a simple example, let  $f$  be a cuspform on  $GL_q$  for  $1 < q < r$ , and define

$$\varphi_{s,f}(p) = \left| \frac{(\det a)^q}{(\det x)^{r-q}} \right|^s \cdot f(x) \quad \left( \text{for } p = \begin{pmatrix} a & b \\ 0 & x \end{pmatrix} \in P^{r-q,q} \right)$$

and extend  $\varphi$  to all of  $GL_r(\mathbb{A})$  by right invariance with respect to the standard maximal compact subgroups  $K_\nu$  at all places  $\nu$  of  $k$ . Form the Eisenstein series

$$E_{s,f}(g) = \sum_{\gamma \in P_k \backslash GL_r(k)} \varphi_{s,f}(g\gamma) \quad (\text{where } P = P^{r-q,q})$$

To normalize this Eisenstein series, let  $\Phi$  be a Schwartz function on  $\mathbb{A}^{q \times q}$ , and suppose that  $\Phi$  is right and left  $K_\nu$ -invariant for all places  $\nu$ . Let  $\xi_f(s)$  be the Godement-Jacquet global zeta integral

$$\xi_f(s) = \int_{GL_q(\mathbb{A})} f(t^{-1}) |\det t|^s \Phi(t) dt$$

We will show that  $\xi_f(rs) \cdot E_{s,f}$  is *entire* in  $s$ .

The argument is a straightforward combination of ideas of [Godement-Jacquet 1972] (using Poisson summation, extending the Tate-Iwasawa argument for  $GL_1$ ) with the Ingham-Rankin-Selberg integral Mellin representation of Eisenstein series (see [Godement 1966] for  $GL_2$ ).

These Eisenstein series are *partly degenerate* in the sense that (except for extreme cases) while including some cuspidal data they do also use non-cuspidal data (determinants) on one of the simple factors of the Levi component. Thus, these Eisenstein series play no direct role in the  $L^2$  spectral theory on the arithmetic quotients, but, rather, are *residues* of the non-degenerate Eisenstein series which *do* enter the spectral theory (as in [Langlands 1976] and [Moeglin Waldspurger 1989], [Moeglin Waldspurger 1995]).

The present special approach gives better results on possible poles of these special semi-degenerate Eisenstein series than an iterated residues viewpoint. This is useful in integral representations of L-functions.

Both the classical treatment of  $GL_2$  and the present integral representation can be viewed as *theta correspondences*, for groups  $GL_n$  viewed as globally split forms of unitary groups.

## 1. Review of a well-known example

First we review some extremely degenerate Eisenstein series (really a *pair* of Eisenstein series), on  $GL_r$  over a number field  $k$ , with continuation properties established by Poisson summation in one step. This material is standard. The (vacuously) *cuspidal data* here is a Hecke character  $\chi$ . The Eisenstein series of this section *may* have poles.

View  $\mathbb{A}^r$  and  $k^r$  as *row* vectors. Let  $e_1, \dots, e_r$  the standard basis for  $k^r$ . We use the  $(r-1, 1)$  parabolic

$$P = \{p \in GL_r : (k \cdot e_r) \cdot p = k \cdot e_r\} = \left\{ \begin{pmatrix} a & * \\ 0 & d \end{pmatrix} : a \in GL_{r-1}, d \in GL_1 \right\} = \text{stabilizer of the line } ke_r$$

We make vectors<sup>[1]</sup>  $\varphi$  constructed from a Hecke character<sup>[2]</sup>  $\chi$  and a Schwartz function  $\Phi$  on  $\mathbb{A}^r$ , by

$$\varphi(g) = \chi(\det g) \int_{\mathbb{J}} \chi(t)^r \Phi(t \cdot e_r \cdot g) dt$$

The  $r^{\text{th}}$  power of  $\chi$  in the integrand and the leading factor of  $\chi(\det g)$  combine to give the invariance  $\varphi(zg) = \varphi(g)$  for all  $z$  in the center  $Z_{\mathbb{A}}$  of  $GL_r(\mathbb{A})$ .

By changing variables in the integral, and by the product formula, we observe the left equivariance

$$\begin{aligned} \varphi(pg) &= \chi(\det pg) \int_{\mathbb{J}} \chi(t)^r \Phi(t \cdot e_r \cdot pg) dt \\ &= \chi(\det a) \chi(d)^{1-r} \cdot \varphi(g) \quad (\text{for } p = \begin{pmatrix} a & b \\ 0 & d \end{pmatrix} \in P_{\mathbb{A}}) \end{aligned}$$

with respect to  $P_{\mathbb{A}}$ .<sup>[3]</sup>

Note that the normalization is *not*  $\varphi(1) = 1$ , but, rather,

$$\varphi(1) = \int_{\mathbb{J}} \chi(t)^r \Phi(t \cdot e_r) dt \quad (\text{Tate-Iwasawa zeta integral at } \chi^r)$$

Denote this zeta integral by  $\xi = \xi(\chi^r, \Phi(0, *))$ , indicating that it only depends upon the values of  $\Phi$  along the last coordinate axis. Thus, we should construe the Eisenstein series, associated to  $\varphi$  in the usual fashion, as having a factor of  $\xi(\chi^r, \Phi(0, *))$  included, namely

$$\xi(\chi^r, \Phi(0, *)) \cdot E(g) = \sum_{\gamma \in P_k \backslash G_k} \varphi(\gamma g) \quad (\text{convergent for } \text{Re}(\chi) \gg 0)$$

[1] These vectors lie in a parametrized family of induced representations, specifically, degenerate principal series, but this fact is not immediately necessary.

[2] As usual, a *Hecke character* is a continuous group homomorphism of  $GL_1(k) \backslash GL_1(\mathbb{A})$  to  $\mathbb{C}^\times$ . The idele class group quotient  $\mathbb{J}/k^\times = GL_1(\mathbb{A})/GL_1(k)$  can be factored as  $(\mathbb{J}^1/k^\times) \times \mathbb{R}^+$ , where  $\mathbb{J}^1$  is ideles with idele norm 1, and  $\mathbb{R}^+$  is the group of positive real numbers, imbedded in  $\mathbb{J}$  by  $t \rightarrow (t^{1/N}, \dots, t^{1/N}, 1, 1, \dots)$ , with  $N = [k : \mathbb{Q}]$ , and  $t^{1/N}$  appears at the archimedean places. A Hecke character  $\chi$  has a corresponding decomposition  $\chi(t) = \chi_1(t) \cdot |t|^s$ , where  $\chi_1$  is a Hecke character *trivial* on the copy of  $\mathbb{R}^+$ , and  $s \in \mathbb{C}$ . As usual, write  $\text{Re}(\chi) = \text{Re}(s)$ .

[3] This left equivariance is exactly the left equivariance required of a vector in *degenerate principal series* representations of  $GL_r(\mathbb{A})$  induced from that character  $p \rightarrow \chi(\det a) \chi(d)^{1-r}$  on  $P_{\mathbb{A}}$ . For now, we have no need of any properties of this representation, but it is convenient that for *generic*  $\chi$  (in the sense of *typical*, rather than in the technical sense of *having a Whittaker model*) these degenerate principal series are *irreducible*. Specifically, these representations are irreducible for the complex parameter  $s$  *off a discrete set of points* (with no limit point in  $\mathbb{C}$ ). Thus, these integral expressions inevitably produce *all* vectors in the representation. This irreducibility is non-trivial, and in the immediate sequel we will not use either it or the implied surjectivity.

Poisson summation proves the meromorphic continuation of this Eisenstein series, just as the Tate-Iwasawa version of Riemann's argument for the Euler-Riemann zeta. Let

$$\mathbb{J}^+ = \{t \in \mathbb{J} : |t| \geq 1\} \quad \mathbb{J}^- = \{t \in \mathbb{J} : |t| \leq 1\}$$

and  $g^t = (g^\top)^{-1}$  (transpose inverse)

$$\begin{aligned} \xi \cdot E(g) &= \sum_{\gamma \in P_k \backslash G_k} \varphi(\gamma g) = \chi(\det g) \sum_{\gamma \in P_k \backslash G_k} \int_{\mathbb{J}} \chi(t)^r \Phi(t \cdot x \cdot \gamma g) dt \\ &= \chi(\det g) \sum_{\gamma \in P_k \backslash G_k} \int_{k^\times \backslash \mathbb{J}} \chi(t)^r \sum_{\lambda \in k^\times} \Phi(t \cdot \lambda e_r \cdot \gamma g) dt \\ &= \chi(\det g) \int_{k^\times \backslash \mathbb{J}} \chi(t)^r \sum_{x \in k^r - 0} \Phi(t \cdot x \cdot g) dt \end{aligned}$$

Let<sup>[4]</sup>

$$\Theta(g) = \sum_{x \in k^r} \Phi(t \cdot x \cdot g)$$

Then

$$\xi \cdot E(g) = \chi(\det g) \int_{k^\times \backslash \mathbb{J}^+} \chi(t)^r [\Theta(g) - \Phi(0)] dt + \chi(\det g) \int_{k^\times \backslash \mathbb{J}^-} \chi(t)^r [\Theta(g) - \Phi(0)] dt = (\text{entire}) + (?)$$

since the usual sort of estimate shows that the integral over  $k^\times \backslash \mathbb{J}^+$  converges absolutely for all  $\chi$ . Following Riemann *et alia*, via Poisson summation we rewrite the second part of the integral as an analogous integral over  $k^\times \backslash \mathbb{J}^+$ . Poisson summation asserts that

$$\sum_{x \in k^r - 0} \Phi(t \cdot x \cdot g) + \Phi(0) = |t|^{-r} |\det g|^{-1} \sum_{x \in k^r - 0} \widehat{\Phi}(t^{-1} \cdot x \cdot g^t) + |t|^{-r} |\det g|^{-1} \widehat{\Phi}(0)$$

using standard change-of-variables properties of the Fourier transform.<sup>[5]</sup> Let

$$\Theta'(g^t) = \sum_{x \in k^r} \widehat{\Phi}(t \cdot x \cdot g^t).$$

Then, removing the  $\Phi(0)$  and  $\widehat{\Phi}(0)$  terms, and replacing  $t$  by  $t^{-1}$  in the integral over  $k^\times \backslash \mathbb{J}^-$  turns this integral into

$$\begin{aligned} &|\det g|^{-1} \chi(\det g) \int_{k^\times \backslash \mathbb{J}^+} (|t| |\chi(t)^{-1}|)^r [\Theta'(g^t) - \widehat{\Phi}(0)] dt \\ &- \chi(\det g) \Phi(0) \int_{k^\times \backslash \mathbb{J}^-} \chi(t)^r dt + |\det g|^{-1} \chi(\det g) \widehat{\Phi}(0) \int_{k^\times \backslash \mathbb{J}^-} \chi(t)^r |t|^{-r} dt \end{aligned}$$

The integral over  $k^\times \backslash \mathbb{J}^+$  is entire. Thus, the non-elementary part of the integral is converted into two *entire* integrals over  $k^\times \backslash \mathbb{J}^+$  together with two elementary integrals that give the only possible poles:

$$\xi \cdot E(g) = (\text{entire}) - \chi(\det g) \Phi(0) \int_{k^\times \backslash \mathbb{J}^-} \chi(t)^r dt + |\det g|^{-1} \chi(\det g) \widehat{\Phi}(0) \int_{k^\times \backslash \mathbb{J}^-} \chi(t)^r |t|^{-r} dt$$

[4] As in Iwasawa-Tate theory, this function  $\theta$  has a role analogous to the Jacobi theta function in the classical argument.

[5] Any non-trivial additive character  $\psi$  on  $\mathbb{A}/k$ , and any non-degenerate pairing  $\langle \cdot, \cdot \rangle$  on  $\mathbb{A}^r$  ( $k$ -valued on  $k^r \times k^r$ ) can be used to define a Fourier transform. The specific choice does not matter, but of course the notion of *transpose* is defined via the pairing.

As in Tate-Iwasawa theory, with  $\chi$  decomposed as

$$\chi(t) = \chi_1(t) \cdot |t|^s \quad (\text{with } \chi_1 \text{ trivial on the copy of } \mathbb{R}^+ \text{ in } \mathbb{J})$$

the relatively elementary integrals can be evaluated

$$\int_{k^\times \backslash \mathbb{J}^-} |t|^{rs} \chi_1(t)^r dt = \int_{k^\times \backslash \mathbb{J}^1} \chi_1(t)^r dt \cdot \int_0^1 t^{rs} dt = \frac{1}{rs} \int_{k^\times \backslash \mathbb{J}^1} \chi_1(t)^r dt = \begin{cases} \frac{\text{vol}(k^\times \backslash \mathbb{J}^1)}{rs} & (\text{for } \chi_1^r = 1) \\ 0 & (\text{for } \chi_1^r \neq 1) \end{cases}$$

Similarly,

$$\int_{k^\times \backslash \mathbb{J}^-} |t|^{r(s-1)} \chi_1(t)^r dt = \begin{cases} \frac{\text{vol}(k^\times \backslash \mathbb{J}^1)}{r(s-1)} & (\text{for } \chi_1^r = 1) \\ 0 & (\text{for } \chi_1^r \neq 1) \end{cases}$$

That is, the only possible poles of  $\xi \cdot E(g)$  are at  $s = 0, 1$ , and these occur only when  $\chi_1^r = 1$ , and the residues are  $\Phi(0)$  at  $s = 0$  and  $\widehat{\Phi}(0)$  at  $s = 1$  (times the volume constant):

$$\begin{aligned} [1.1] \quad \xi(\chi^r, \Phi(0, *)) \cdot E(g) &= \chi(\det g) \int_{k^\times \backslash \mathbb{J}^+} \chi(t)^r [\Theta(g) - \Phi(0)] dt \\ &\quad + |*| \chi^{-1}(\det g^t) \int_{k^\times \backslash \mathbb{J}^+} |*| \chi^{-1}(t)^r [\Theta'(g^t) - \widehat{\Phi}(0)] dt \\ &\quad + \begin{cases} -\frac{\chi(\det g) \text{vol}(k^\times \backslash \mathbb{J}^1)}{rs} + \frac{|*| \chi^{-1}(\det g^t) \text{vol}(k^\times \backslash \mathbb{J}^1)}{r(s-1)} & \text{for } \chi^r = 1 \\ 0 & \text{for } \chi^r \neq 1 \end{cases} \end{aligned}$$

Thus, the possible poles of  $E(g)$  itself, *without* the zeta factor  $\xi(\chi^r, \Phi(0, *))$ , may be at  $s = 0, 1$  and at the zeros of  $\xi$ . For reasonable choices of  $\Phi$ , these zeros occur only in the strip  $0 < \text{Re}(s) < \frac{1}{r}$ .

[1.2] **Remark:** The estimate here on possible poles is stronger than the otherwise-natural estimates that arise when these Eisenstein series are analyzed as residues of non-degenerate Eisenstein series.

[1.3] **Remark:** Although it is a good indicator of analytic properties, the above expression for  $E(g)$  requires further interpretation to give a functional equation connecting it to another Eisenstein series. This is explicated in the next section.

## 2. Simplest non-self-associate functional equation

Except for the case  $r = 2$ , the functional equation of the degenerate Eisenstein series (above) attached to the  $(r-1, 1)$  parabolic does *not* relate it to itself. We will see that there is a functional equation relating it to an Eisenstein series for the  $(1, r-1)$  parabolic [6]

$$Q = \left\{ \begin{pmatrix} 1\text{-by-1} & * \\ 0 & (r-1)\text{-by-}(r-1) \end{pmatrix} \right\}$$

Indeed, the discussion of meromorphic continuation presents  $g^t$  inside  $\widehat{\Phi}$ . In the special case  $r = 2$ , up to a scalar,  $g^t$  is *conjugate* to  $g$ .

[6] The fact that the diagonal block sizes in  $P$  and  $Q$  are merely permuted is no coincidence. In a broader context, we would say that these two parabolics are *associate*.

Let  $E^P = E^P(\chi, \Phi)$  be the Eisenstein series of the previous section. We want to understand the Fourier-transformed part of the integral expression for  $E^P$  as being made from a left-equivariant function  $\varphi'$  which should be roughly like  $\varphi$ . An obvious candidate,

$$\varphi'(g) = \chi(g^t) \int_{\mathbb{J}} \chi(t)^r \widehat{\Phi}(t \cdot e_r \cdot g^t) dt \quad (?)$$

is left equivariant by the image  $P^\iota$  of  $P$  under the involution  $\iota$ , but  $P^\iota$  is not a standard parabolic subgroup. Further, except for  $r = 2$ , the subgroup  $P^\iota$  is not even *conjugate* to  $P$ , but is conjugate to the (standard) parabolic  $Q$  via

$$w_\circ P^\iota w_\circ = Q \quad (\text{where } w_\circ = \begin{pmatrix} & & & 1 \\ & & \cdot & \\ & & & \\ 1 & & & \end{pmatrix}, \text{ anti-diagonal})$$

Thus, unsurprisingly, noting that  $e_1 = e_r \cdot w_\circ$ , we should take

$$\varphi'(g) = \chi(g^\sigma) \int_{\mathbb{J}} \chi(t)^r \widehat{\Phi}(t \cdot e_1 \cdot g^t) dt = \text{left equivariant by } Q_{\mathbb{A}} \quad (\text{note occurrence of } e_1, \text{ not } e_r)$$

As before, the normalization involves a zeta integral

$$\xi(\chi^r, \widehat{\Phi}(*, 0)) = \varphi'(1) = \int_{\mathbb{J}} \chi(t)^r \widehat{\Phi}(t \cdot e_1) dt$$

Thus, we should take the viewpoint that this zeta integral occurs as an implied factor when we form the Eisenstein series  $E^Q(\chi^r, \widehat{\Phi})$  attached to  $\varphi'$ , namely

$$\xi(\chi^r, \widehat{\Phi}(*, 0)) \cdot E^Q(g) = \sum_{\gamma \in Q_k \backslash GL_r(k)} \varphi'(\gamma g)$$

The visible symmetry in the integral expression for  $E^P$  from above gives the functional equation

$$\xi(\chi^r, \widehat{\Phi}(*, 0)) \cdot E^P(\chi, \Phi) = \xi((|*|\chi^{-1})^{-1}, \widehat{\Phi}(0, *)) \cdot E^Q(|*|\chi^{-1}, \widehat{\Phi})$$

### 3. Half-cuspidal data

Take  $1 < q < r$ , and consider the standard maximal proper parabolic

$$P = P^{r-q, q} = \left\{ \begin{pmatrix} (r-q)\text{-by-}(r-q) & * \\ 0 & q\text{-by-}q \end{pmatrix} \right\}$$

We will demonstrate the meromorphic continuation of an Eisenstein series attached to *degenerate* data on the copy of  $GL_{r-q}$ , and to *cuspidal* data on the copy of  $GL_r$ . The general form of the discussion is parallel to the previous.

Let  $f$  be a cuspform on  $GL_q(\mathbb{A})$ , in the strong sense that  $f$  is in  $L^2(GL_q(k) \backslash GL_q(\mathbb{A})^1)$ ,  $f$  meets the Gelfand-Fomin condition<sup>[7]</sup>

$$\int_{N_k \backslash N_{\mathbb{A}}} f(n g) dn = 0 \quad (\text{for almost all } g)$$

[7] In words, this condition is *integrating to 0 over horocycles*.

and  $f$  generates an irreducible representation of  $GL_q(k_\nu)$  locally at all places  $\nu$  of  $k$ . For a Schwartz function  $\Phi$  on  $\mathbb{A}^{q \times r}$  and Hecke character  $\chi$ , let

$$\varphi(g) = \varphi_{\chi, f, \Phi}(g) = \chi(\det g)^q \int_{GL_q(\mathbb{A})} f(h^{-1}) \chi(\det h)^r \Phi(h \cdot [0_{q \times (r-q)} \ 1_q] \cdot g) dh$$

This function  $\varphi$  has the same central character as  $f$ . It is left invariant by the adèle points of the unipotent radical

$$N = \left\{ \begin{pmatrix} 1_{r-q} & * \\ 0 & 1_r \end{pmatrix} \right\} \quad (\text{unipotent radical of } P = P^{r-q, q})$$

The function  $\varphi$  is left invariant under the  $k$ -rational points of the standard Levi component of  $P$ ,

$$M = \left\{ \begin{pmatrix} a & 0 \\ 0 & d \end{pmatrix} : a \in GL_{r-q}, d \in GL_r \right\}$$

To understand the normalization, observe that

$$\xi(\chi^r, f, \Phi(0, *)) = \varphi(1) = \int_{GL_q(\mathbb{A})} f(h^{-1}) \chi(\det h)^r \Phi(h \cdot [0_{q \times (r-q)} \ 1_q]) dh$$

is a Godement-Jacquet zeta integral for the standard L-function attached to the cusform  $f$  (or perhaps a contragredient). Thus, the corresponding Eisenstein series includes this zeta integral as a factor, so write <sup>[8]</sup>

$$\xi(\chi^r, f, \Phi(0, *)) \cdot E_{\chi, f, \Phi}^P(g) = \sum_{\gamma \in P_k \backslash GL_r(k)} \varphi(\gamma g) \quad (\text{convergent for } \operatorname{Re}(\chi) \gg 0)$$

Now prove the meromorphic continuation via Poisson summation.

$$\begin{aligned} & \xi(\chi^r, f, \Phi(0, *)) \cdot E_{\chi, f, \Phi}^P(g) \\ &= \chi(\det g)^q \sum_{\gamma \in P_k \backslash GL_r(k)} \int_{GL_q(k) \backslash GL_q(\mathbb{A})} f(h) \chi(\det h)^{-r} \sum_{\alpha \in GL_q(k)} \Phi(h^{-1} \cdot [0 \ \alpha] \cdot g) dh \\ &= \chi(\det g)^q \int_{GL_q(k) \backslash GL_q(\mathbb{A})} f(h) \chi(\det h)^{-r} \sum_{y \in k^q \times r, \text{ full rank}} \Phi(h^{-1} \cdot y \cdot g) dh \end{aligned}$$

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[8] To understand  $\varphi$ , observe that for  $p = \begin{pmatrix} a & b \\ 0 & x \end{pmatrix}$  in  $P_{\mathbb{A}}$ ,

$$\begin{aligned} \varphi(p) &= \chi(\det p)^q \int_{GL_q(\mathbb{A})} f(h^{-1}) \chi(\det h)^r \Phi(h \cdot [0_{q \times (r-q)} \ x]) dh \\ &= \chi(\det p)^q \chi(\det x)^{-r} \int_{GL_q(\mathbb{A})} f(xh) \chi(\det h^{-1})^r \Phi(h^{-1} \cdot [0_{q \times (r-q)} \ 1_q]) dh \end{aligned}$$

by replacing  $h$  by  $h^{-1}$  then replacing  $h$  by  $xh$ . For simplicity, suppose that  $f$  is a *spherical* vector everywhere locally,

$$\int_{GL_q(\mathbb{A})} f(xh) \chi(\det h^{-1})^r \Phi(h^{-1} \cdot [0_{q \times (r-q)} \ 1_q]) dh = f(x) \cdot \int_{GL_q(\mathbb{A})} f(h) \chi(\det h^{-1})^r \Phi(h^{-1} \cdot [0_{q \times (r-q)} \ 1_q]) dh$$

and the last integral is a Godement-Jacquet zeta integral  $\xi(\chi^r, f, \Phi(0, *))$ .

The Gelfand-Fomin condition on  $f$  will compensate for the otherwise-irksome full-rank constraint. Anticipating that we can drop the rank condition suggests that we define

$$\Theta_\Phi(h, g) = \sum_{y \in k^{q \times r}} \Phi(h^{-1} \cdot y \cdot g)$$

Now we show (as in Godement-Jacquet) that the non-full-rank terms integrate to 0. [9]

**[3.1] Proposition:** For  $f$  a cuspform, less-than-full-rank terms integrate to 0, that is,

$$\int_{GL_q(k) \backslash GL_q(\mathbb{A})} f(h) \chi(\det h)^{-r} \sum_{y \in k^{q \times r}, \text{rank} < q} \Phi(h^{-1} \cdot y \cdot g) dh = 0$$

*Proof:* Since this is asserted for arbitrary Schwartz functions  $\Phi$ , we can take  $g = 1$ . By linear algebra, given  $y_0 \in k^{q \times r}$  of rank  $\ell$ , there is  $\alpha \in GL_q(k)$  such that

$$\alpha \cdot y_0 = \begin{pmatrix} \text{rank-}\ell \text{ } \ell\text{-by-}r \text{ block} \\ 0_{(q-\ell) \times r} \end{pmatrix}$$

Thus, without loss of generality fix  $y_0$  of the latter shape. Let  $Y$  be the orbit of  $y_0$  under left multiplication by the rational points of the parabolic

$$P^{\ell, q-\ell} = \begin{pmatrix} \ell\text{-by-}\ell & * \\ 0 & (q-\ell)\text{-by-}(q-\ell) \end{pmatrix} \subset GL_q$$

This is some set of matrices of the same shape as  $y_0$ . Then the subsum over  $GL_q(k) \cdot y_0$  is

$$\int_{GL_q(k) \backslash GL_q(\mathbb{A})} f(h) \chi(\det h)^{-r} \sum_{y \in GL_q(k) \cdot y_0} \Phi(h^{-1} \cdot y) dh = \int_{P_k^{\ell, q-\ell} \backslash GL_q(\mathbb{A})} f(h) \chi(\det h)^{-r} \sum_{y \in Y} \Phi(h^{-1} \cdot y) dh$$

Let  $N$  and  $M$  be the unipotent radical and standard Levi component of  $P^{\ell, q-\ell}$ , namely, usual,

$$N = \begin{pmatrix} 1_\ell & * \\ 0 & 1_{q-\ell} \end{pmatrix} \quad M = \begin{pmatrix} \ell\text{-by-}\ell & 0 \\ 0 & (q-\ell)\text{-by-}(q-\ell) \end{pmatrix}$$

Then the integral can be rewritten as an iterated integral

[9] There are issues of convergence. First, for  $\text{Re}(\chi)$  sufficiently large, the integral

$$\chi(\det g)^q \sum_{\gamma \in P_k \backslash GL_r(k)} \int_{GL_q(k) \backslash GL_q(\mathbb{A})} f(h) \chi(\det h)^{-r} \Theta(h, g) dh$$

is absolutely convergent. Also, we have the integrals analogous to integrals over  $k^\times \backslash \mathbb{J}^+$ . That is, let  $GL_q^+ = \{h \in GL_q(\mathbb{A}) : |\det h| \geq 1\}$ . Then, for *arbitrary*  $\chi$ , using the fact that cuspforms  $f$  are of rapid decay in Siegel sets,

$$\chi(\det g)^q \sum_{\gamma \in P_k \backslash GL_r(k)} \int_{GL_q(k) \backslash GL_q^+(\mathbb{A})} f(h) \chi(\det h)^{-r} \sum_{y \in k^{q \times r}, \text{full rank}} \Phi(h^{-1} \cdot y \cdot g) dh$$

is absolutely convergent.

$$\begin{aligned}
& \int_{N_k M_k \backslash GL_q(\mathbb{A})} f(h) \chi(\det h)^{-r} \sum_{y \in Y} \Phi(h^{-1} \cdot y) dh \\
&= \int_{N_{\mathbb{A}} M_k \backslash GL_q(\mathbb{A})} \sum_{y \in Y} \int_{N_k \backslash N_{\mathbb{A}}} f(nh) \chi(\det nh)^{-r} \Phi((nh)^{-1} \cdot y) dn dh \\
&= \int_{N_{\mathbb{A}} M_k \backslash GL_q(\mathbb{A})} \sum_{y \in Y} \chi(\det h)^{-r} \Phi(h^{-1} \cdot y) \left( \int_{N_k \backslash N_{\mathbb{A}}} f(nh) dn \right) dh
\end{aligned}$$

since all fragments but  $f(nh)$  in the integrand are left invariant by  $N_{\mathbb{A}}$ . But the inner integral of  $f(nh)$  is 0, by the Gelfand-Fomin condition, so the whole is 0. ///

Let  $\iota$  denote the transpose-inverse involution(s). Poisson summation gives

$$\begin{aligned}
\Theta_{\Phi}(h, g) &= \sum_{y \in k^{q \times r}} \Phi(h^{-1} \cdot y \cdot g) \\
&= |\det(h^{-1})^{\iota}|^r |\det g^{\iota}|^q \sum_{y \in k^{q \times r}} \widehat{\Phi}((h^{\iota})^{-1} \cdot y \cdot g^{\iota}) = |\det(h^{-1})^{\iota}|^r |\det g^{\iota}|^q \Theta_{\widehat{\Phi}}(h^{\iota}, g^{\iota})
\end{aligned}$$

As with  $\Theta_{\Phi}$ , the not-full-rank summands in  $\Theta_{\widehat{\Phi}}$  integrate to 0 against cuspforms. Thus, letting

$$GL_q^+ = \{h \in GL_q(\mathbb{A}) : |\det h| \geq 1\} \quad GL_q^- = \{h \in GL_q(\mathbb{A}) : |\det h| \leq 1\}$$

$$\begin{aligned}
& \xi(\chi^r, f, \Phi(0, *)) \cdot E_{\chi, f, \Phi}^P(g) = \chi(\det g)^q \int_{GL_q(k) \backslash GL_q(\mathbb{A})} f(h) \chi(\det h)^{-r} \Theta_{\Phi}(h, g) dh \\
&= \chi(\det g)^q \int_{GL_q(k) \backslash GL_q^+} f(h) \chi(\det h)^{-r} \Theta_{\Phi}(h, g) dh + \chi(\det g)^q \int_{GL_q(k) \backslash GL_q^-} f(h) \chi(\det h)^{-r} \Theta_{\Phi}(h, g) dh \\
&= \chi(\det g)^q \int_{GL_q(k) \backslash GL_q^+} f(h) \chi(\det h)^{-r} \Theta_{\Phi}(h, g) dh \\
&\quad + \chi(\det g)^q \int_{GL_q(k) \backslash GL_q^-} |\det(h^{-1})^{\iota}|^r |\det g^{\iota}|^q f(h) \chi(\det h)^{-r} \Theta_{\widehat{\Phi}}(h^{\iota}, g^{\iota}) dh
\end{aligned}$$

By replacing  $h$  by  $h^{\iota}$  in the second integral we convert it to an integral over  $GL_q(k) \backslash GL_q^+$ , and the whole is

$$\begin{aligned}
& \xi(\chi^r, f, \Phi(0, *)) \cdot E_{\chi, f, \Phi}^P(g) = \chi(\det g)^q \int_{GL_q(k) \backslash GL_q^+} f(h) \chi(\det h)^{-r} \Theta_{\Phi}(h, g) dh \\
&\quad + |*| \chi^{-1}(\det g^{\iota})^q \int_{GL_q(k) \backslash GL_q^+} f(h^{\iota}) |*| \chi^{-1}(\det h^{\iota})^{-r} \Theta_{\widehat{\Phi}}(h, g^{\iota}) dh
\end{aligned}$$

Since  $f \circ \iota$  is a cuspform, the second integral is entire in  $\chi$ . Thus, we have proven

$$\xi(\chi^r, f, \Phi(0, *)) \cdot E_{\chi, f, \Phi}^P \text{ is entire}$$

**[3.2] Remark:** Except for the extreme case  $q = r - 1$ , these Eisenstein series are degenerate, so occur only as residues of (purely) cuspidal-data Eisenstein series. Assessing poles of residues seems less effective in the present special circumstances than the above argument.



## 4. Associate Eisenstein series

The functional equation of the Eisenstein series  $E_{\chi, f, \Phi}^P$  relates it not to itself but to an Eisenstein series for the parabolic  $Q$  associate to  $P$ , namely<sup>[10]</sup>

$$Q = P^{q, r-q} = \left\{ \begin{pmatrix} q\text{-by-}q & * \\ 0 & (r-q)\text{-by-}(r-q) \end{pmatrix} \right\}$$

The functions

$$\varphi'(g) = \chi(\det g^t)^q \int_{GL_q(\mathbb{A})} f(h^t) \chi(\det h)^{-r} \widehat{\Phi}(h^{-1} \cdot [1_q \ 0_{q \times (r-q)}] \cdot g^t) dh$$

arising in the meromorphic continuation for  $E_{\chi, f, \Phi}^P$  are vectors in induced representations for  $Q$ . The values  $\varphi'(1)$  are Godement-Jacquet zeta integrals

$$\xi(\chi^r, f \circ \iota, \widehat{\Phi}(*0)) = \int_{GL_q(\mathbb{A})} f(h^t) \chi(\det h)^{-r} \widehat{\Phi}(h^{-1} \cdot [1_q \ 0_{q \times (r-q)}]) dh$$

Then form an Eisenstein series by

$$\xi(\chi^r, f \circ \iota, \widehat{\Phi}(*0)) \cdot E_{\chi, f \circ \iota, \widehat{\Phi}}^Q = \sum_{\gamma \in Q_k \backslash GL_r(k)} \varphi'(\gamma g)$$

The expression above that proves the meromorphic continuation also has a visible symmetry, giving the functional equation

$$\xi(\chi^r, f, \Phi(0*)) \cdot E_{\chi, f, \Phi}^P = \xi((|\cdot| \chi^{-1})^r, f \circ \iota, \widehat{\Phi}(*0)) \cdot E_{|\chi^{-1}, f \circ \iota, \widehat{\Phi}}^Q$$

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<sup>[10]</sup> This applies even in the symmetrical case  $r - q = q$ , since the roles of the cuspidal data and degenerate data are reversed in the functional equation.