





Math and Climate Seminar IMA




Mathematics and Climate Research Network


**Joint MCRN/IMA Math and Climate Seminar**  
 Tuesdays 11:15 – 12:05  
 streaming video available at  
 www.ima.umn.edu

**Budyko's Model as a Dynamical System**  
 Richard McGehee



Seminar on the Mathematics of Climate  
 IMA, MCRN, School of Mathematics  
 November 22, 2013




**Budyko's Model**  
 Budyko's Equation

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y)) - (A + BT) + C(\bar{T} - T)$$

Labels: surface temperature, heat capacity, insolation, albedo, OLR, heat transport,  $\bar{T} = \int_0^1 T(y) dy$ , sin(latitude)

Symmetry assumption:  $0 \leq y = \sin(\text{latitude}) \leq 1$

Chylek and Coakley's quadratic approximation:  
 $s(y) \approx 1 - 0.241(3y^2 - 1)$



**Budyko's Model**  
 Budyko's Equation


$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y)) - (A + BT) + C(\bar{T} - T)$$

equilibrium solution:  $T = T^*(y)$

$$Qs(y)(1 - \alpha(y)) - (A + BT^*(y)) + C(\bar{T} - T^*(y)) = 0$$

Integrate:  
 $\int_0^1 (Qs(y)(1 - \alpha(y)) - (A + BT^*(y)) + C(\bar{T} - T^*(y))) dy = 0$   
 $Q(1 - \bar{\alpha}) - (A + B\bar{T}^*) = 0$   
 where  $\bar{\alpha} = \int_0^1 \alpha(y)s(y) dy$

Global mean temperature at equilibrium  
 $\bar{T}^* = \frac{1}{B}(Q(1 - \bar{\alpha}) - A)$



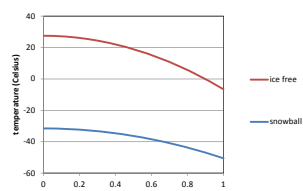

**Budyko's Model**  
 Budyko's Equilibrium

$$Qs(y)(1 - \alpha(y)) - (A + BT^*(y)) + C(\bar{T}^* - T^*(y)) = 0$$

Global mean temperature at equilibrium:  
 $\bar{T}^* = \frac{1}{B}(Q(1 - \bar{\alpha}) - A)$  ( $\bar{\alpha} = \int_0^1 \alpha(y)s(y) dy$ )

Equilibrium temperature profile:  
 $T^*(y) = \frac{1}{B+C}(Qs(y)(1 - \alpha(y)) - A + C\bar{T}^*)$

Example:  
 $C = 3.04$   
 $\alpha = 0.32$ : ice free  
 $\alpha = 0.62$ : snowball

**Budyko's Model**  
 Ice Albedo Feedback

What if the albedo is not constant?

**Ice Line Assumption:** There is a single ice line at  $y = \eta$  between the equator and the pole. The albedo is  $\alpha_1$  below the ice line and  $\alpha_2$  above it.

$$\alpha(y, \eta) = \begin{cases} \alpha_1 & y < \eta \\ \alpha_2 & y > \eta \end{cases}$$

Equilibrium condition:  
 $Qs(y)(1 - \alpha(y, \eta)) - (A + BT^*(y)) + C(\bar{T}^* - T^*(y)) = 0$

Equilibrium solution:  
 $T^*(y) = \frac{1}{B+C}(Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*)$   
 where  $\bar{T}^* = \frac{1}{B}(Q(1 - \bar{\alpha}(\eta)) - A)$  ( $\bar{\alpha}(\eta) = \int_0^1 \alpha(y, \eta)s(y) dy$ )

**Budyko's Model**

**Ice Albedo Feedback**

$$T_e^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y,\eta)) - A + CT_e^*)$$

For each fixed  $\eta$ , there is an equilibrium solution for Budyko's equation.

**Budyko's Model**

**Dynamics**

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y)) - (A+BT) + C(\bar{T} - T)$$

Labels: surface temperature, heat capacity, insolation, albedo, OLR, heat transport,  $\bar{T} = \int_0^1 T(y) dy$

Let  $X$  be the space of functions where  $T$  lives. (e.g.  $L^1([0,1])$ )  
 Let  $L: X \rightarrow X: LT = C\bar{T} - (B+C)T$ ,  
 $f(y) = Qs(y)(1-\alpha(y)) - A$

Budyko's equation can be written as a linear vector field on  $X$ .

$$R \frac{dT}{dt} = f + LT$$

**Budyko's Model**

**Budyko's Equation**

$$R \frac{dT}{dt} = f + LT$$

$$LT(y) = C \int_0^1 T(\xi) d\xi - (B+C)T(y)$$

If  $X$  is a decent space (e.g.  $L^1([0,1])$ ), then  $L$  is a continuous linear operator on  $X$ .

Equilibrium solution:  
 $T^* = -L^{-1}f$   
 if  $L$  is invertible.

Stability depends on the spectrum of  $L$ .

**Budyko's Model**

**Linear Operator**

$$LT(y) = C \int_0^1 T(\xi) d\xi - (B+C)T(y)$$

Let  $X_0 = \{T \in X: \int_0^1 T(y) dy = 0\}$ ,  $X_1 = \{T \in X: T \text{ is constant}\}$

Then  $X = X_0 \oplus X_1$   
 $LT = -(B+C)T$ , for  $T \in X_0$   
 $LT = -BT$ , for  $T \in X_1$

**L is invertible.**

**Spectrum**

- $-B$  is an eigenvalue, with eigenspace  $X_1$  (dimension 1)
- $-(B+C)$  is an eigenvalue, with eigenspace  $X_0$  (codimension 1)

**The equilibrium solution is stable.**

**Budyko's Model**

**Ice Albedo Feedback**

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y,\eta)) - (A+BT) + C(\bar{T} - T)$$

For each fixed  $\eta$ , there is a **stable** equilibrium solution for Budyko's equation.

How to pick one?

**Budyko's Model**

**Ice Albedo Feedback**

For each fixed  $\eta$ , there is a stable equilibrium solution for Budyko's equation.

**Standard assumption:** Permanent ice forms if the annual average temperature is below  $T_c = -10^\circ\text{C}$  and melts if the annual average temperature is above  $T_c$ .

**Additional condition:** The average temperature across the ice boundary is the critical temperature  $T_c$ .

$$\frac{1}{2}(T_e^*(\eta+) + T_e^*(\eta-)) = T_c = -10$$



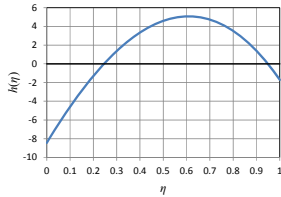
### Budyko's Model Ice Albedo Feedback

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

The additional condition:  $\frac{1}{2}(T_e^*(\eta^+) + T_e^*(\eta^-)) = T_c = -10$   
can be written:

$$h(\eta) \equiv \frac{Q}{B+C} \left( s(\eta)(1 - \alpha_0) + \frac{C}{B} \left[ 1 - \alpha_2 + (\alpha_2 - \alpha_1) \int_0^\eta s(y) dy \right] \right) - \frac{A}{B} - T_c = 0$$

Two equilibria (zeros of  $h$ ) satisfy the additional condition.

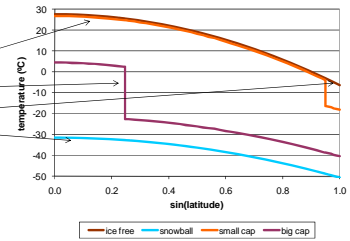


### Budyko's Model Ice Albedo Feedback

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

Interesting Solutions:

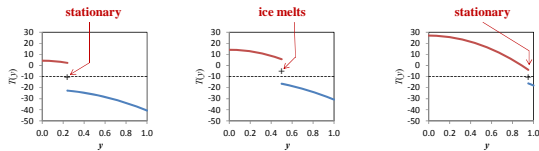
- small cap
- large cap
- ice free
- snowball



### Budyko's Model Dynamics of the Ice Line

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

**Idea:**  
If the average temperature across the ice line is above the critical temperature, some ice will melt, moving the ice line toward the pole. If it is below the critical temperature, the ice will advance toward the equator.



Widiasih's equation:  $\frac{d\eta}{dt} = \varepsilon(T(\eta) - T_c)$



### Budyko's Model Dynamics of the Ice Line

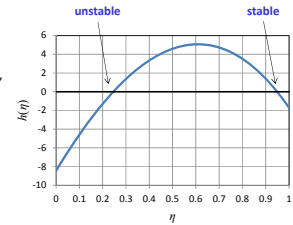
$$\frac{d\eta}{dt} = \varepsilon(T(\eta) - T_c)$$

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

State space:  $[0, 1] \times X$

**Widiasih's Theorem.** For sufficiently small  $\varepsilon$ , the system has an attracting invariant curve given by the graph of a function  $\Phi_\varepsilon : [0, 1] \rightarrow X$ . On this curve, the dynamics are approximated by the equation

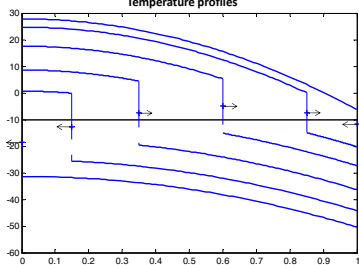
$$\frac{d\eta}{dt} = \varepsilon h(\eta)$$



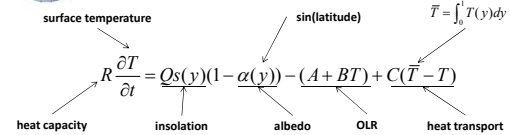
### Budyko's Model Budyko-Widiasih Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta)$$

Temperature profiles



### Budyko's Model Summary



reduces to

$$\frac{d\eta}{dt} = \varepsilon h(\eta)$$

$$\frac{d\eta}{dt} = \varepsilon h(\eta) = \varepsilon \left( \frac{Q}{B+C} \left( s(\eta)(1 - \alpha_0) + \frac{C}{B} \left[ 1 - \alpha_2 + (\alpha_2 - \alpha_1) \int_0^\eta s(y) dy \right] \right) - \frac{A}{B} - T_c \right)$$

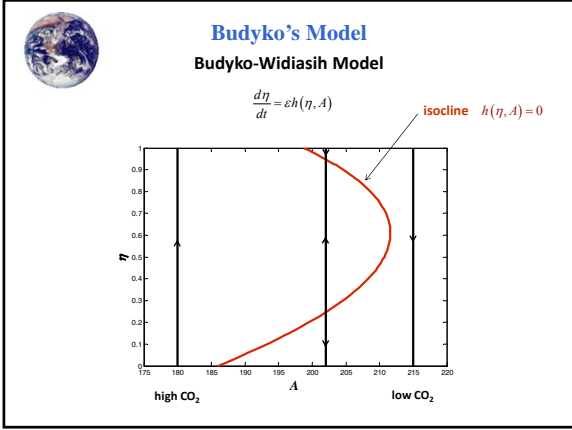
**Budyko's Model**  
Budyko-Widiasih Model

$$R \frac{\partial T}{\partial t} = Q_s(y)(1 - \alpha(y)) - (A + BT) + C(\bar{T} - T)$$

Labels: surface temperature,  $\sin(\text{latitude})$ , heat capacity, insolation, albedo, OLR, heat transport,  $\bar{T} = \int_0^1 T(y) dy$

**What about the greenhouse effect?**

$A + BT$  is the outgoing long wave radiation. This term decreases if the greenhouse gases increase. We view  $A$  as a parameter.

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A)$$


**Budyko's Model**  
Budyko-Widiasih Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A)$$

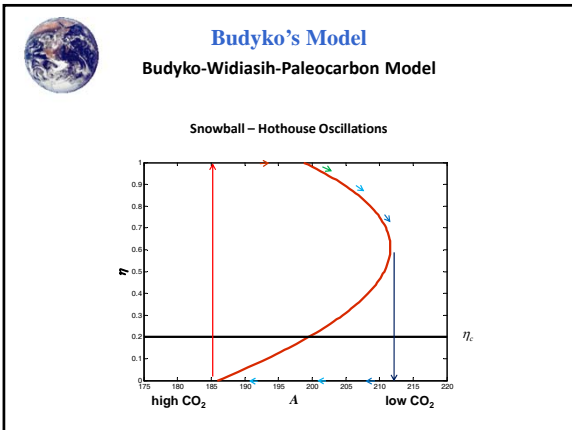
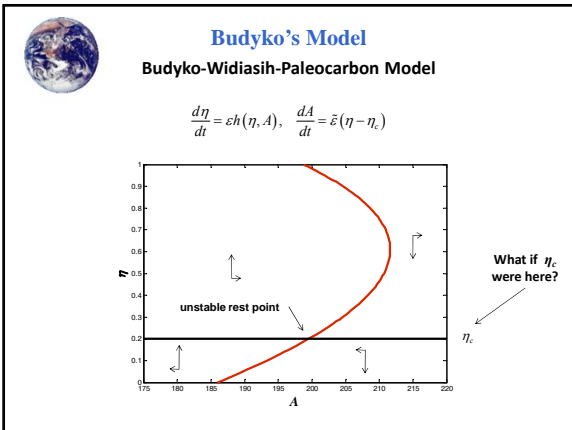
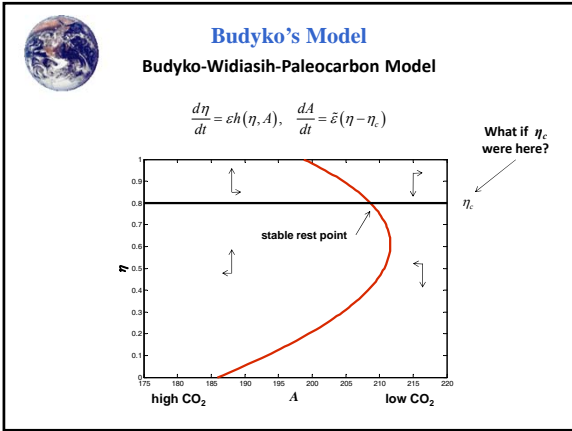
What if  $A$  is a dynamical variable?

Simple equation:

$$\frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c) \quad \leftarrow \text{MCRN Paleocarbon equation}$$

$0 < \tilde{\varepsilon} \ll \varepsilon \ll 1$

New system:

$$\begin{cases} \frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c) \\ \frac{d\eta}{dt} = \varepsilon h(\eta, A) \end{cases}$$




### Budyko's Model

#### Water Belt Model

The snowball Earth hypothesis is not universally accepted. Complex ocean-dwelling photosynthetic organisms appear to have survived through the time when the snowball would have occurred. Perhaps the equator never fully froze over, leaving a belt of open water covering the equator. Hence "water belt" planet, aka "Jormungand".

D.S. Abbot, A. Voigt, and D. Koll (2011), The Jormungand global climate state and implications for Neoproterozoic glaciations, *J. Geophys. Res.* **116**.



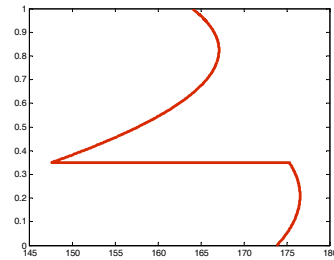
### Budyko's Model

#### Water Belt Model

$h$  is modified, so the isocline looks like this:

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A),$$

$$\frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c)$$

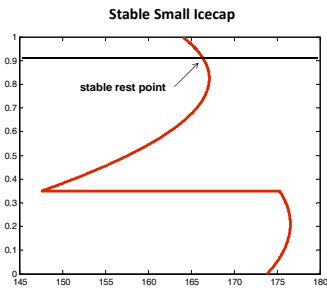


### Budyko's Model

#### Water Belt Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A),$$

$$\frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c)$$

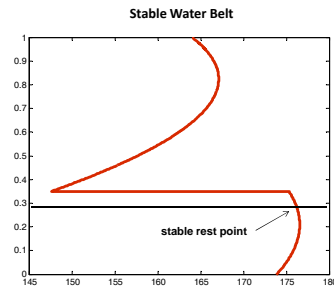


### Budyko's Model

#### Water Belt Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A),$$

$$\frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c)$$

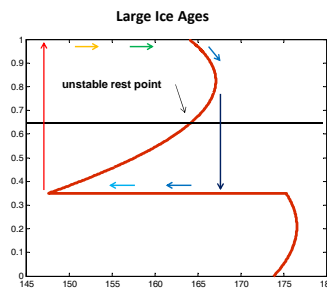


### Budyko's Model

#### Water Belt Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A),$$

$$\frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c)$$

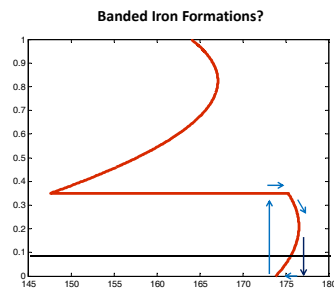


### Budyko's Model

#### Water Belt Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A),$$

$$\frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c)$$





## Budyko's Model

### Budyko's Equation

$$R \frac{\partial T}{\partial t} = Q_s(y)(1 - \alpha(y)) - (A + BT) + C(\bar{T} - T)$$

Labels for the equation components:

- surface temperature:  $T$
- heat capacity:  $R$
- insolation:  $Q_s(y)$
- albedo:  $\alpha(y)$
- OLR:  $A + BT$
- heat transport:  $C(\bar{T} - T)$
- $\bar{T} = \int_0^1 T(y) dy$
- sin(latitude):  $y$

Next Time  
Glacial Cycles