

Dynamics of Energy Balance Models for Planetary Climate

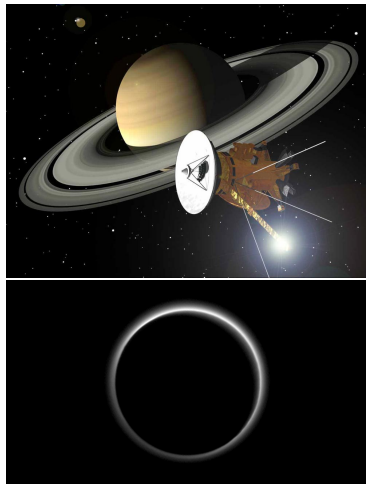
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Motivation

- Low dimensional climate models are important for understanding the predominant forces affecting the climate of Earth
- Fly-bys of Pluto and other rocky celestial bodies in our solar system have raised interest in other climates



Photos from nasa.gov

Energy Balance Models in 1969

- Budyko and Sellers (independently) proposed energy balance models for the Earth (1, 14)
- Wanted to study if another glacial age was possible
- Both models had the same major components: incoming solar radiation, outgoing radiation, and energy transfer:



$$R\Delta T = Q(y)(1 - \alpha(y)) - (A + BT) + \Gamma(T)$$

Energy Balance Models Today

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + \Gamma(T)$$

where

$$\Gamma(T(y, \eta)) = -C \left(T(y, \eta) - \int_0^1 T(\gamma, \eta) d\gamma \right).$$

Widiasih introduced an equation for the dynamics of the ice line, η , in 2012 (18)

$$\frac{d}{dt} \eta = \epsilon(T(y, \eta) - T_c)$$

“Quadratic Approximation”

In (9) McGehee and Widiasih consider the Budyko-type equation for $y, \eta \in [0, 1]$

$$\frac{\partial}{\partial t} T = \frac{1}{R} (Q_s(y)(1 - \alpha(\eta, y)) - (A + BT(y, \eta)) - C(T(y, \eta) - \bar{T}))$$

with dynamic ice line

$$\dot{\eta} = \rho(T(y, \eta) - T_c)$$

and piecewise constant albedo function

$$\alpha(y, \eta) = \begin{cases} \alpha_w & y < \eta \\ \alpha_0 & y = \eta \\ \alpha_i & y > \eta \end{cases}$$

where

$$\alpha_0 = \frac{\alpha_w + \alpha_i}{2}.$$

Piecewise Function Space

The equilibrium temperature profile can be written

$$T_{\eta}^*(y) = \begin{cases} \frac{1}{B+C} (Qs(y)(1 - \alpha_w) - A + C\overline{T_{\eta}^*}) & y < \eta \\ \frac{1}{B+C} (Qs(\eta)(1 - \alpha_0) - A + C\overline{T_{\eta}^*}) & y = \eta \\ \frac{1}{B+C} (Qs(y)(1 - \alpha_i) - A + C\overline{T_{\eta}^*}) & y > \eta \end{cases}$$

which motivates the four-dimensional function space X whose elements are of the form

$$T(y) = \begin{cases} w_0 + \frac{1}{2}z_0 + (w_2 + \frac{1}{2}z_2) p_2(y) & y < \eta \\ w_0 + w_2 p_2(\eta) & y = \eta \\ w_0 - \frac{1}{2}z_0 + (w_2 - \frac{1}{2}z_2) p_2(y) & y > \eta \end{cases}$$

New Budyko's "Equation"

Reformulating the $\partial_t T$ equation in this function space gives

$$R\dot{w}_0 = Q(1 - \alpha_0) - A - Bw_0 + C \left(\left(\eta - \frac{1}{2} \right) z_0 + z_2 \int_0^\eta p_2(y) dy \right)$$

$$R\dot{z}_0 = Q(\alpha_i - \alpha_w) - (B + C)z_0$$

$$R\dot{w}_2 = Qs_2(1 - \alpha_0) - (B + C)w_2$$

$$R\dot{z}_2 = Qs_2(\alpha_i - \alpha_w) - (B + C)z_2$$

and the ice line equation becomes

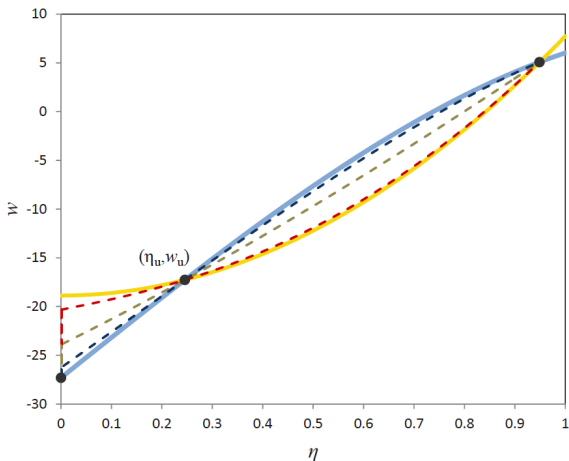
$$R\dot{\eta} = \rho \left(w_0 - \frac{Q(1 - \alpha_0)}{B + C} s_2 p_2(\eta) + T_c \right).$$

Write

$$R\dot{w}_0 = -B(w_0 - F(\eta))$$

$$R\dot{\eta} = -\rho(w_0 - G(\eta))$$

(η, w_0) Phase Space



$$R\dot{w}_0 = -B(w_0 - F(\eta))$$

$$R\dot{\eta} = -\rho(w_0 - G(\eta))$$

The Jormungand Model

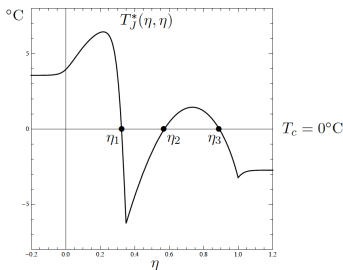
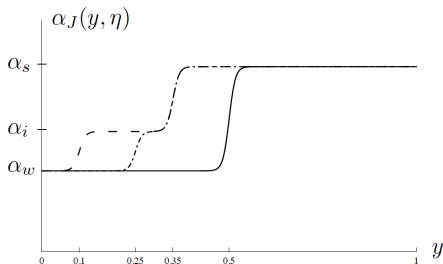
Define the function

$$\delta(\eta) = \begin{cases} -\eta + .35, & \eta < .35 \\ 0, & \eta \geq .35 \end{cases}$$

which represents the extent of the bare ice and the Jormungand albedo function

$$\alpha_J(y, \eta) = \frac{\alpha_s + \alpha_w}{2} + \frac{\alpha_i - \alpha_w}{2} \tanh(M(y - \eta)) \\ + \frac{\alpha_s - \alpha_i}{2} \tanh(M(y - (\eta + \delta(\eta))))$$

Jormungand albedo function and Temperature Profile



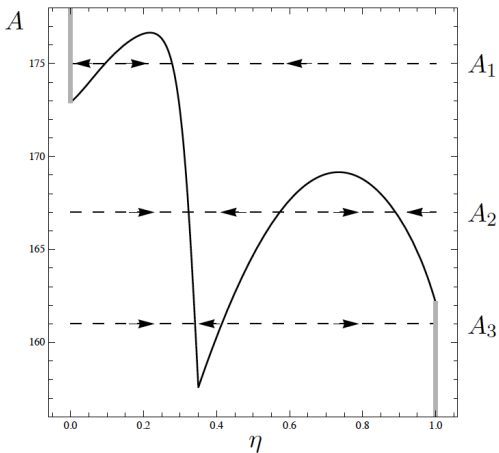
- Widiasih's Theorem still applies with this albedo function
- There is a locally attracting invariant manifold

$$\mathcal{P}_J^* = \{(\Phi_J^*(\eta), \eta) : \eta \in \mathbb{R}\}$$

within $\mathcal{O}(\epsilon)$ of the manifold of fixed points

$$\mathcal{T}_J^* = \{(T_J^*(y, \eta), \eta) : \eta \in \mathbb{R}\}$$

Jormungand Bifurcation in the Greenhouse Gas Parameter



Insolation Distribution on Rapidly Spinning Planets

- Actual distribution can be found using orbital parameters (as seen in (5, 8, 17)):

$$s(y, \beta) = \frac{2}{\pi^2} \int_0^{2\pi} \sqrt{1 - \left(\sqrt{1 - y^2} \sin \beta \sin \gamma - y \cos \beta \right)^2} d\gamma$$

- We use Legendre approximations in EBMs because the above integral doesn't have a closed form expression. Instead use

$$s(y, \beta) = \sum_{n=0}^{\infty} s_{2n}(\beta) p_{2n}(y)$$

Insolation Distribution on Rapidly Spinning Planets

Write the $2n$ -th degree Legendre polynomial as

$$p_{2n}(y) = \sum_{k=0}^n a_{2k} y^{2k}$$

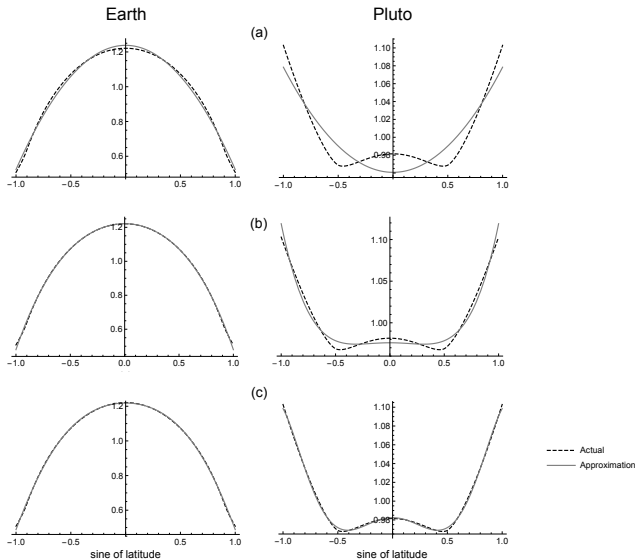
and

$$s_{2n}(\beta) = P_{2n} \sum_{k=0}^n a_{2k} c_{2k}(\beta).$$

where

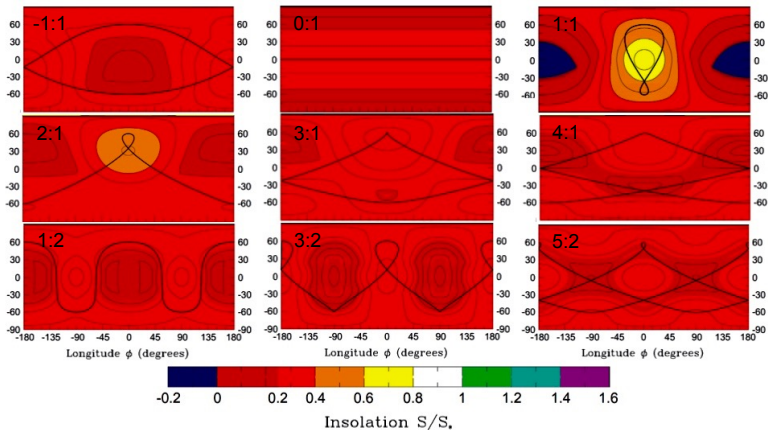
$$c_{2k}(\beta) = \sum_{j=0}^k \binom{2k}{2j} \frac{(\sin \beta)^{2(k-j)} (\cos \beta)^{2j}}{\pi^2} \left(\int_{-\pi/2}^{\pi/2} (\cos \hat{\phi})^{2(k+1-j)} (\sin \hat{\phi})^{2j} d\hat{\phi} \right) \left(\int_0^{2\pi} (\cos \hat{\theta})^{2(k-j)} d\hat{\theta} \right)$$

Insolation Distribution on Rapidly Spinning Planets



Small Integer Spin-Orbit Resonances

Mean annual insolation distributions for obliquity $\beta = 60^\circ$ and eccentricity $e = .2$:



A. Dobrovolskis, "Insolation on exoplanets with eccentricity and obliquity."

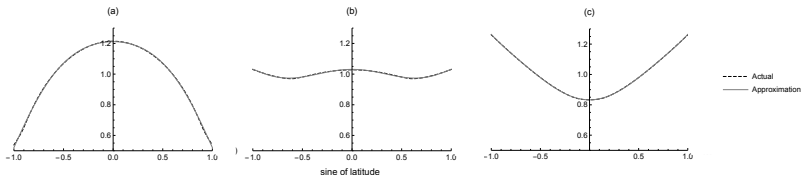
Open Questions about Insolation

- Can we quantify when we can use the “rapidly spinning planet” method/formula and still have error less than τ in our approximation?

- Can we find a closed form expression for insolation on planets with small integer resonances?

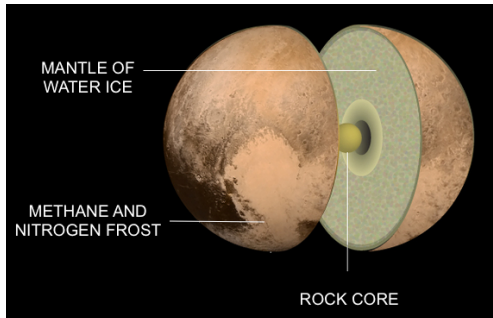
Open Questions about Insolation

- For a reasonable range of parameter values, the Budyko map for Pluto doesn't have any nontrivial stable fixed points. Why?
 - Could Pluto's "upside down" insolation be playing a factor?
 - Is it because Pluto's insolation is relatively flat?



Open Questions about Ice Lines

- Reformulate model to accommodate ice planets
 - Is more than one ice present and how do we account for different albedos?
 - How are the ices situated on the surface? Are ices mixed? Are they layered?



From nasa.gov

Open Questions about Ice Lines

- Remove symmetry assumptions?
 - Investigate four ice lines $(\eta_{SP}, \eta_{SE}, \eta_{NE}, \eta_{NP})$ with the properties
 - (i) $\eta_{SP}, \eta_{SE}, \eta_{NE}, \eta_{NP} \in [-1, 1]$.
 - (ii) $-1 \leq \eta_{SP} \leq \eta_{SE} \leq \eta_{NE} \leq \eta_{NP} \leq 1$.
 - (iii) Ice is located between η_{SP} and η_{SE} and between η_{NE} and η_{NP} .
 - Initial investigations into this case show potential oscillations in the ice line dynamics

Other Open Questions Concerning EBMs

- How do we accurately model a planet whose atmosphere “freezes out” (as Pluto’s might)?
- Is the diffusion model more accurate for a planet with no oceans? For a planet without an atmosphere?
 - In which cases do we get the same results with both models?
 - In which cases do we get different results?

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