



**An Introduction to Budyko's Model**

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Mathematics of Climate Seminar  
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**Budyko's Model**  
Dynamical Models


Add Heat Transport  $R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha) - (A+BT) + C(\bar{T} - T)$

global mean temperature  $\bar{T}(t) = \int_0^1 T(y,t) dy$

Second Law of Thermodynamics:  
Energy travels from hot places to cold places.

*Equilibrium temperature profile?*

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**Budyko's Model**


Conservation of Energy

temperature change  $\sim$  energy in - energy out

short wave energy from the Sun  $\rightarrow$   $\rightarrow$  temperature change  
 long wave energy from the Earth  $\leftarrow$   $\leftarrow$  temperature change

*Everything else is detail.*

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**Budyko's Model**  
Budyko's Equation


$R \frac{\partial T}{\partial t} = \underbrace{Qs(y)}_{\text{insolation}} (1 - \underbrace{\alpha(y)}_{\text{albedo}}) - (A + \underbrace{BT}_{\text{OLR}}) + C(\underbrace{\bar{T}}_{\text{heat transport}} - T)$

Labels: surface temperature, sin(latitude), heat capacity, insolation, albedo, OLR, heat transport

Symmetry assumption:  $0 \leq y = \sin(\text{latitude}) \leq 1$

Chylek and Coakley's quadratic approximation:  
 $s(y) \approx 1 - 0.241(3y^2 - 1)$


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**Budyko's Model**  
Dynamical Models

	Model	Equilibrium
Perfectly Thermally Conducting Black Body	$R \frac{dT}{dt} = Q - \sigma T^4$	$T = (Q/\sigma)^{1/4}$
Plus Albedo	$R \frac{dT}{dt} = Q(1-\alpha) - \sigma T^4$	$T = ((1-\alpha)Q/\sigma)^{1/4}$
Switch to Surface Temperature	$R \frac{dT}{dt} = Q(1-\alpha) - (A+BT)$	$T = ((1-\alpha)Q - A)/B$
Dependence on Latitude	$R \frac{\partial T(y,t)}{\partial t} = Qs(y)(1-\alpha) - (A+BT(y,t))$	$T(y) = ((1-\alpha)Qs(y) - A)/B$

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**Budyko's Model**  
Budyko's Equilibrium

$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y)) - (A+BT) + C(\bar{T} - T)$

albedo depends on latitude

equilibrium solution:  $T = T^*(y)$

$Qs(y)(1-\alpha(y)) - (A+BT^*(y)) + C(\bar{T} - T^*(y)) = 0$

Integrate:

$$\int_0^1 (Qs(y)(1-\alpha(y)) - (A+BT^*(y)) + C(\bar{T} - T^*(y))) dy = 0$$

$$Q \int_0^1 s(y) dy - Q \int_0^1 s(y) \alpha(y) dy - A \int_0^1 dy - B \int_0^1 T^*(y) dy + C \left( \int_0^1 \bar{T} dy - \int_0^1 T^*(y) dy \right) = 0$$

$Q(1-\bar{\alpha}) - (A+BT^*) = 0$

equilibrium global mean temperature  $\bar{T}^* = \frac{1}{B}(Q(1-\bar{\alpha}) - A)$

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**Budyko's Model**  
**Budyko's Equilibrium**

$Qs(y)(1-\alpha(y)) - (A + BT^*(y)) + C(\bar{T}^* - T^*(y)) = 0$

Global mean temperature at equilibrium:  
 $\bar{T}^* = \frac{1}{B} (Q(1-\bar{\alpha}) - A)$  ( $\bar{\alpha} = \int_0^1 \alpha(y)s(y)dy$ )

Solve for  $T^*(y)$ .  
 $Qs(y)(1-\alpha(y)) - A + C\bar{T}^* = BT^*(y) + CT^*(y) = (B+C)T^*(y)$   
 $T^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y)) - A + C\bar{T}^*)$

Equilibrium temperature profile:  
 $T^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y)) - A + C\bar{T}^*)$

where  $\bar{T}^* = \frac{1}{B} (Q(1-\bar{\alpha}) - A)$  and  $\bar{\alpha} = \int_0^1 \alpha(y)s(y)dy$

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**Budyko's Model**  
**Ice Albedo Feedback**

temperature warms  
ice melts  
albedo decreases  
more sunlight absorbed  
temperature warms  
REPEAT

Why would it stop?

Warmer temperatures  
Less snow and ice  
More sunlight absorbed by land and sea

<http://www.i-fink.com/melting-polar-ice/>

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**Budyko's Model**  
**Budyko's Equilibrium**

$Qs(y)(1-\alpha(y)) - (A + BT^*(y)) + C(\bar{T}^* - T^*(y)) = 0$

Equilibrium temperature profile:  $T^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y)) - A + C\bar{T}^*)$

$C = 3.04$   
 $\alpha(y) = 0.32$ : ice free  
 $\alpha(y) = 0.62$ : snowball (constant albedo)

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**Budyko's Model**  
**Ice Albedo Feedback**

temperature warms  
ice melts  
albedo decreases  
more sunlight absorbed  
temperature warms  
REPEAT

Why would it stop?

M. I. Budyko, "The effect of solar radiation variations on the climate of the Earth," *Tellus XXI*, 611-619, 1969.

[http://www.inenco.org/index\\_principals.html](http://www.inenco.org/index_principals.html)

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**Budyko's Model**  
**Budyko's Equilibrium**

ice won't melt (no exit from snowball)  
ice will form (icecap)

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**Budyko's Model**  
**Ice Albedo Feedback**

Why would it stop?  
**Budyko's Equation**

surface temperature  $\frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y)) - (A + BT) + C(\bar{T} - T)$

heat capacity  $R$   $\frac{\partial T}{\partial t}$   $Qs(y)$   $1-\alpha(y)$   $A$   $B$   $T$   $C$   $\bar{T}$   $T$


insolation  $Qs(y)$   $1-\alpha(y)$   $A$   $B$   $T$   $C$   $\bar{T}$   $T$

albedo  $1-\alpha(y)$   $A$   $B$   $T$   $C$   $\bar{T}$   $T$

OLR  $A + BT$   $C$   $\bar{T}$   $T$

heat transport  $C(\bar{T} - T)$

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### Budyko's Model

#### Ice Albedo Feedback

What if the albedo is not constant?

**Ice Line Assumption:** There is a single ice line at  $y=\eta$  between the equator and the pole. The albedo is  $\alpha_1$  below the ice line and  $\alpha_2$  above it.

$$\alpha(y, \eta) = \begin{cases} \alpha_1 & y < \eta \\ \alpha_2 & y > \eta \end{cases}$$

Equilibrium condition:


$$Qs(y)(1-\alpha(y, \eta)) - (A + BT_s^*(y)) + C(\bar{T} - T_s^*(y)) = 0$$

Equilibrium solution:

$$T_s^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y, \eta)) - A + C\bar{T})$$

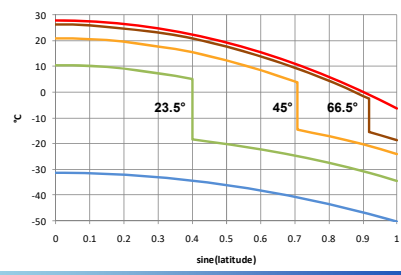
where  $\bar{T} = \frac{1}{B} (Q(1-\bar{\alpha}(\eta)) - A)$  ( $\bar{\alpha}(\eta) = \int_0^1 \alpha(y, \eta) s(y) dy$ )

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### Budyko's Model

#### Ice Albedo Feedback

$$T_s^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y, \eta)) - A + C\bar{T})$$


For each fixed  $\eta$ , there is an equilibrium solution for Budyko's equation

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### Budyko's Model

#### Ice Albedo Feedback

equilibrium temperature profile:

$$T_s^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y, \eta)) - A + C\bar{T}), \text{ where } \bar{T} = \frac{1}{B} (Q(1-\bar{\alpha}(\eta)) - A)$$

albedo:

$$\alpha(y, \eta) = \begin{cases} \alpha_1 = 0.32, & y < \eta \\ \alpha_2 = 0.62, & y > \eta \end{cases}$$

global albedo:

$$\bar{\alpha}(\eta) = \int_0^1 \alpha(y, \eta) s(y) dy = \int_0^\eta \alpha_1 s(y) dy + \int_\eta^1 \alpha_2 s(y) dy$$

let:

$$S(\eta) = \int_0^\eta s(y) dy, \quad 1 - S(\eta) = \int_\eta^1 s(y) dy, \quad \text{since } 1 = \int_0^1 s(y) dy$$


then:

$$\bar{\alpha}(\eta) = \alpha_1 S(\eta) + \alpha_2 (1 - S(\eta)) = \alpha_2 - (\alpha_2 - \alpha_1) S(\eta) = 0.62 - 0.3S(\eta)$$

$$S(\eta) = \int_0^\eta s(y) dy = \int_0^\eta (1 - 0.241(3y^2 - 1)) dy = \eta - 0.241(\eta^3 - \eta)$$

Chylek & Coakley

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### Budyko's Model

#### Dynamics

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

Let  $X$  be the space of functions where  $T$  lives. (e.g.  $L^1([0,1])$ )  
Let  $L: X \rightarrow X: LT = C\bar{T} - (B+C)T$ ,  
 $f(y) = Qs(y)(1-\alpha(y, \eta)) - A$


Budyko's equation can be written as a linear vector field on  $X$ .

$$R \frac{dT}{dt} = f + LT$$

The operator  $L$  has only point spectrum, with all eigenvalues negative.  
Therefore, all solutions are stable.  
True for any albedo function.

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### Budyko's Model

#### Ice Albedo Feedback

equilibrium temperature profile:

$$T_s^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y, \eta)) - A + C\bar{T}), \text{ where } \bar{T} = \frac{1}{B} (Q(1-\bar{\alpha}(\eta)) - A)$$

albedo:

$$\alpha(y, \eta) = \begin{cases} \alpha_1 = 0.32, & y < \eta \\ \alpha_2 = 0.62, & y > \eta \end{cases}$$

global albedo:

$$\bar{\alpha}(\eta) = \int_0^1 \alpha(y, \eta) s(y) dy = \int_0^\eta \alpha_1 s(y) dy + \int_\eta^1 \alpha_2 s(y) dy$$

let:

$$S(\eta) = \int_0^\eta s(y) dy, \quad 1 - S(\eta) = \int_\eta^1 s(y) dy, \quad \text{since } 1 = \int_0^1 s(y) dy$$


then:

$$\bar{\alpha}(\eta) = \alpha_1 S(\eta) + \alpha_2 (1 - S(\eta)) = \alpha_2 - (\alpha_2 - \alpha_1) S(\eta) = 0.62 - 0.3S(\eta)$$

$$S(\eta) = \int_0^\eta s(y) dy = \int_0^\eta (1 - 0.241(3y^2 - 1)) dy = \eta - 0.241(\eta^3 - \eta)$$

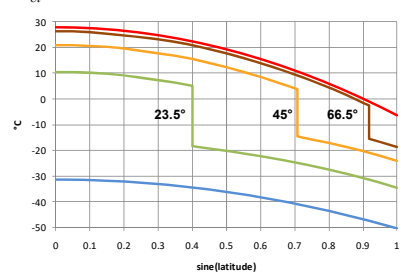
Chylek & Coakley

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### Budyko's Model


#### Ice Albedo Feedback

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$


For each fixed  $\eta$ , there is a **globally stable** equilibrium solution for Budyko's equation.

**How to pick one?**

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### Budyko's Model

#### Ice Albedo Feedback


**Summary**

If we artificially hold the ice line at a fixed latitude, then the surface temperature will come to an equilibrium.

However, if the temperature is high, we would expect ice to melt and the ice line to retreat to higher latitudes. If the temperature is low, we would expect ice to form and the ice line to advance to lower latitudes.

*How to model this expectation?*

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### Budyko's Model

#### Ice Albedo Feedback

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

Equilibrium:  $T_{\eta}^*(y) = \frac{1}{B+C} (Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}_{\eta}^*)$

Ice line condition:  $\frac{1}{2} (T_{\eta}^*(\eta+) + T_{\eta}^*(\eta-)) = T_c = -10$


Albedo:  $\alpha(y, \eta) = \begin{cases} \alpha_1 = 0.32, & y < \eta \\ \alpha_2 = 0.62, & y > \eta \end{cases} \quad \alpha(\eta-, \eta) = \alpha_1, \quad \alpha(\eta+, \eta) = \alpha_2$

$$T_{\eta}^*(\eta+) = \frac{1}{B+C} (Qs(\eta)(1 - \alpha_2) - A + C\bar{T}_{\eta}^*) \quad T_{\eta}^*(\eta-) = \frac{1}{B+C} (Qs(\eta)(1 - \alpha_1) - A + C\bar{T}_{\eta}^*)$$

Ice line condition:  $\frac{1}{2} (T_{\eta}^*(\eta+) + T_{\eta}^*(\eta-)) = \frac{1}{B+C} (Qs(\eta)(1 - \alpha_0) - A + C\bar{T}_{\eta}^*) = T_c = -10$

where:  $\alpha_0 = \frac{\alpha_1 + \alpha_2}{2} = 0.47$

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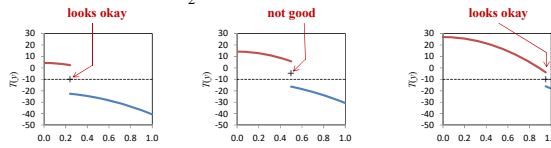
### Budyko's Model

#### Ice Albedo Feedback


For each fixed  $\eta$ , there is a stable equilibrium solution for Budyko's equation.

**Standard assumption:** Permanent ice forms if the annual average temperature is below  $T_c = -10^\circ\text{C}$  and melts if the annual average temperature is above  $T_c$ .

**Additional condition:** The average temperature across the ice boundary is the critical temperature  $T_c$ .

$$\frac{1}{2} (T_{\eta}^*(\eta+) + T_{\eta}^*(\eta-)) = T_c = -10$$


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### Budyko's Model

#### Ice Albedo Feedback

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

Ice line condition:  $\frac{1}{B+C} (Qs(\eta)(1 - \alpha_0) - A + C\bar{T}_{\eta}^*) = T_c = -10$

Rewrite:  $h(\eta) = \frac{1}{B+C} (Qs(\eta)(1 - \alpha_0) - A + C\bar{T}_{\eta}^*) - T_c = 0$


Recall equilibrium GMT:  $\bar{T}_{\eta}^* = \frac{1}{B} (Q(1 - \bar{\alpha}(\eta)) - A)$

Recall average albedo:  $\bar{\alpha}(\eta) = \int_0^{\eta} \alpha(y, \eta) s(y) dy = \alpha_2 - (\alpha_2 - \alpha_1) S(\eta) = 0.62 - 0.3S(\eta)$

where:  $S(\eta) = \int_0^{\eta} s(y) dy = \eta - 0.241(\eta^3 - \eta)$

$$h(\eta) \equiv \frac{Q}{B+C} \left( s(\eta)(1 - \alpha_0) + \frac{C}{B} (1 - \alpha_2 + (\alpha_2 - \alpha_1) S(\eta)) \right) - \frac{A}{B} - T_c = 0$$

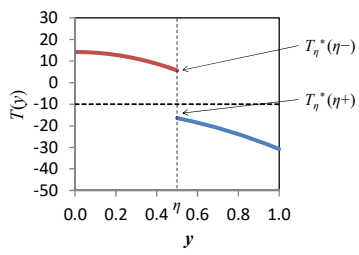
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### Budyko's Model

#### Ice Albedo Feedback

ice line condition:  $\frac{1}{2} (T_{\eta}^*(\eta+) + T_{\eta}^*(\eta-)) = T_c = -10$



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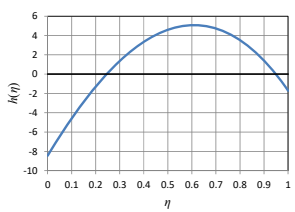
### Budyko's Model

#### Ice Albedo Feedback

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

The additional condition:  $\frac{1}{2} (T_{\eta}^*(\eta+) + T_{\eta}^*(\eta-)) = T_c = -10$

can be written:  $h(\eta) \equiv \frac{Q}{B+C} \left( s(\eta)(1 - \alpha_0) + \frac{C}{B} (1 - \alpha_2 + (\alpha_2 - \alpha_1) S(\eta)) \right) - \frac{A}{B} - T_c = 0$



Two equilibria (zeros of  $h$ ) satisfy the additional condition.

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### Budyko's Model

#### Ice Albedo Feedback

Equilibrium temperature profiles

$$T_{\eta}^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y,\eta)) - A + C\bar{T}_{\eta}^*)$$

Interesting Solutions:

- small cap
- large cap
- ice free
- snowball

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### Budyko's Model

#### Budyko-Widiasih Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta)$$

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### Budyko's Model

#### Dynamics of the Ice Line

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y,\eta)) - (A+BT) + C(\bar{T}-T)$$

Idea:

If the average temperature across the ice line is above the critical temperature, some ice will melt, moving the ice line toward the pole. If it is below the critical temperature, the ice will advance toward the equator.

Widiasih's equation:

$$\frac{d\eta}{dt} = \varepsilon (T(\eta) - T_c)$$

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### Budyko's Model

#### Summary

surface temperature  $\bar{T} = \int_0^1 T(y) dy$

$$R \frac{\partial T}{\partial t} = \underbrace{Qs(y)}_{\text{insolation}} (1 - \underbrace{\alpha(y)}_{\text{albedo}}) - \underbrace{(A+BT)}_{\text{OLR}} + \underbrace{C(\bar{T}-T)}_{\text{heat transport}}$$

heat capacity

reduces to

$$\frac{d\eta}{dt} = \varepsilon h(\eta) = \varepsilon \left( \frac{Q}{B+C} (s(\eta)(1-\alpha_0) + \frac{C}{B}(1-\alpha_2 + (\alpha_2 - \alpha_1)S(\eta))) - \frac{A}{B} - T_c \right)$$

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### Budyko's Model

#### Dynamics of the Ice Line

$$\frac{d\eta}{dt} = \varepsilon (T(\eta) - T_c)$$

State space:  $[0, 1] \times X$

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y,\eta)) - (A+BT) + C(\bar{T}-T)$$

**Widiasih's Theorem.** For sufficiently small  $\varepsilon$ , the system has an attracting invariant curve given by the graph of a function  $\Phi_{\varepsilon}: [0, 1] \rightarrow X$ . On this curve, the dynamics are approximated by the equation

$$\frac{d\eta}{dt} = \varepsilon h(\eta)$$

experts only

Esther R. Widiasih, Dynamics of the Budyko Energy Balance Model, *SIAM J. Appl. Dyn. Syst.*, 12(4), 2068-2092.

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### Budyko's Model

#### Budyko-Widiasih Model

surface temperature  $\bar{T} = \int_0^1 T(y) dy$

$$R \frac{\partial T}{\partial t} = \underbrace{Qs(y)}_{\text{insolation}} (1 - \underbrace{\alpha(y)}_{\text{albedo}}) - \underbrace{(A+BT)}_{\text{OLR}} + \underbrace{C(\bar{T}-T)}_{\text{heat transport}}$$

heat capacity

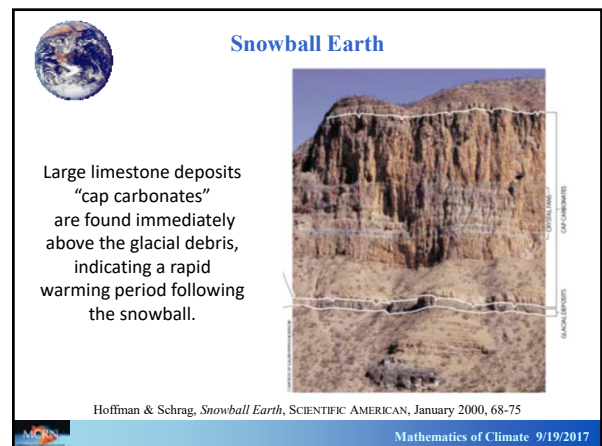
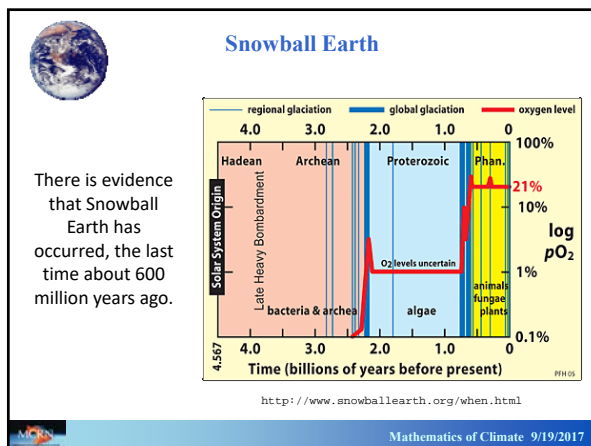
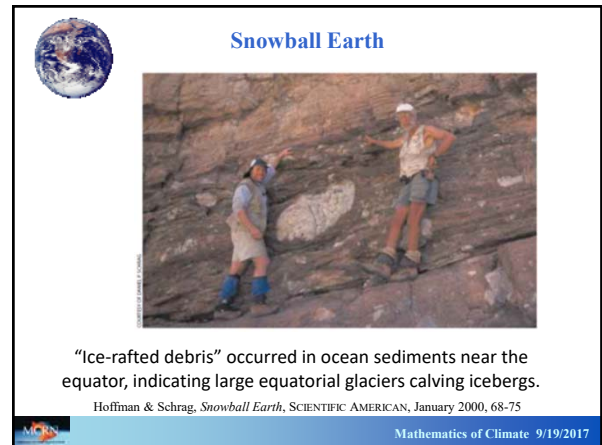
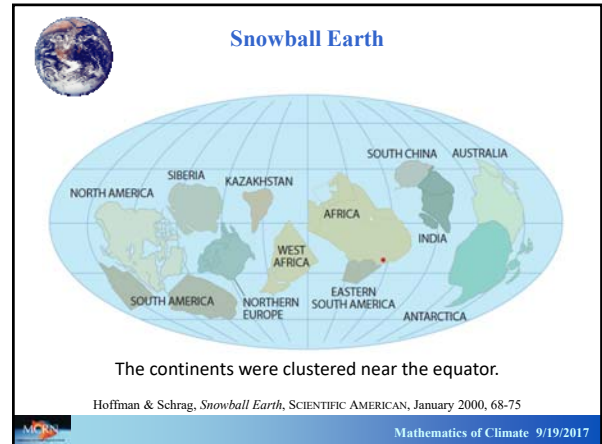
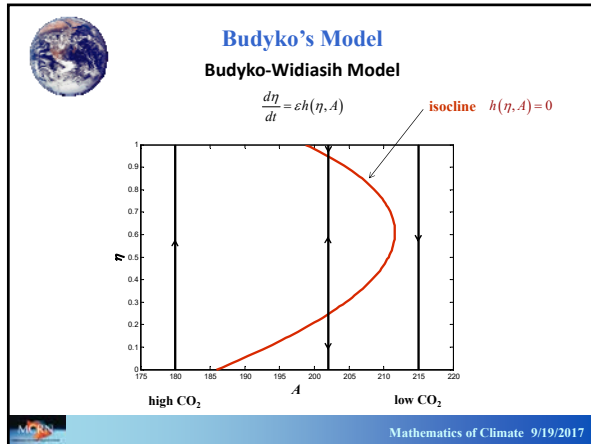
**What about the greenhouse effect?**


$A+BT$  is the outgoing long wave radiation. This term decreases if the greenhouse gases increase.

We view  $A$  as a parameter.

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A)$$

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### Snowball Earth

**Idea:**

When the Earth is mostly ice-covered, silicate weathering slows down, but volcanic activity stays the same, allowing for a build-up of CO2 in the atmosphere.

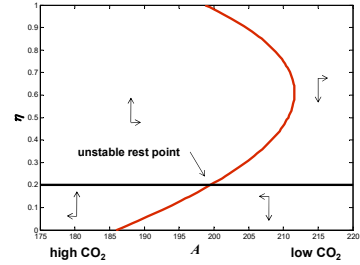
When the Earth is mostly ice-free, silicate weathering speeds up, drawing down the CO2 in the atmosphere.

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### Budyko's Model

#### Budyko-Widiasih-Paleocarbon Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A), \quad \frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c)$$


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### Budyko's Model

#### Budyko-Widiasih Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A)$$

What if  $A$  is a dynamical variable?

**Simple equation:**


$$\frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c)$$

$0 < \tilde{\varepsilon} \ll \varepsilon \ll 1$  ← MCRN Paleocarbon equation (silicate weathering)

**New system:**

$$\begin{cases} \frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c) \\ \frac{d\eta}{dt} = \varepsilon h(\eta, A) \end{cases}$$

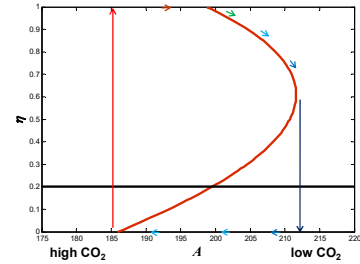
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### Budyko's Model

#### Budyko-Widiasih-Paleocarbon Model

##### Snowball – Hothouse Oscillations

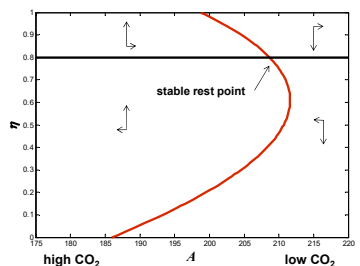


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### Budyko's Model

#### Budyko-Widiasih-Paleocarbon Model


$$\frac{d\eta}{dt} = \varepsilon h(\eta, A), \quad \frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c)$$


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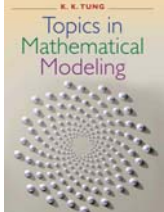


### Budyko's Model

#### Suggested Reading



Hoffman & Schrag, *Snowball Earth*, SCIENTIFIC AMERICAN, January 2000, 68-75



K.K. Tung, *Topics in Mathematical Modeling*, PRINCETON UNIVERSITY PRESS, 2007, Chapter 8

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