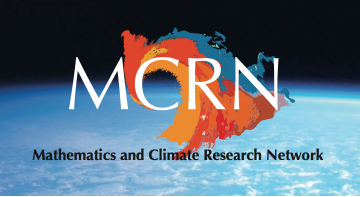



An Introduction to Budyko's Energy Balance Model

Richard McGehee
School of Mathematics
University of Minnesota

Mathematics of Climate Seminar
October 1, 2019

Budyko's Model

Conservation of Energy


temperature change ~ energy in - energy out

short wave energy
from the Sun

long wave energy
from the Earth

Everything else is detail.

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Budyko's Model

Dynamical Models

	Model	Equilibrium
Perfectly Thermally Conducting Black Body	$R \frac{dT}{dt} = Q - \sigma T^4$	$T = (Q/\sigma)^{1/4}$
Plus Albedo	$R \frac{dT}{dt} = Q(1-\alpha) - \sigma T^4$	$T = ((1-\alpha)Q/\sigma)^{1/4}$
Switch to Surface Temperature	$R \frac{dT}{dt} = Q(1-\alpha) - (A + BT)$	$T = ((1-\alpha)Q - A)/B$
Dependence on Latitude	$R \frac{\partial T(y,t)}{\partial t} = Qs(y)(1-\alpha) - (A + BT(y,t))$	$T(y) = ((1-\alpha)Qs(y) - A)/B$

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Budyko's Model

Dynamical Models


Add Heat Transport $R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha) - (A + BT) + C(\bar{T} - T)$

global mean temperature $\bar{T}(t) = \int_0^1 T(y,t) dy$

Second Law of Thermodynamics:
Energy travels from hot places to cold places.

Equilibrium temperature profile?

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Budyko's Model

Budyko's Equation

surface temperature

heat capacity

$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y)) - (A + BT) + C(\bar{T} - T)$

insolation

albedo

OLR

heat transport


sin(latitude)

$\bar{T} = \int_0^1 T(y) dy$

Symmetry assumption: $0 \leq y = \sin(\text{latitude}) \leq 1$

Chylek and Coakley's quadratic approximation:
 $s(y) \approx 1 - 0.241(3y^2 - 1)$

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Budyko's Model

Budyko's Equilibrium

$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y)) - (A + BT) + C(\bar{T} - T)$

albedo depends on latitude

equilibrium solution: $T = T^*(y)$

$Qs(y)(1-\alpha(y)) - (A + BT^*(y)) + C(\bar{T} - T^*(y)) = 0$

Integrate:

$$\int_0^1 (Qs(y)(1-\alpha(y)) - (A + BT^*(y)) + C(\bar{T} - T^*(y))) dy = 0$$

$$Q \int_0^1 s(y) dy - Q \int_0^1 s(y)\alpha(y) dy - A \int_0^1 dy - B \int_0^1 T^*(y) dy + C \left(\int_0^1 \bar{T} dy - \int_0^1 T^*(y) dy \right) = 0$$

$Q(1-\bar{\alpha}) - (A + B\bar{T}^*) = 0$

equilibrium global mean temperature $\bar{T}^* = \frac{1}{B}(Q(1-\bar{\alpha}) - A)$

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Budyko's Model

Budyko's Equilibrium

$$Qs(y)(1-\alpha(y)) - (A + BT^*(y)) + C(\bar{T} - T^*(y)) = 0$$

Solve for $T^*(y)$.

Global mean temperature at equilibrium:

$$\bar{T} = \frac{1}{B} (Q(1-\bar{\alpha}) - A) \quad \left(\bar{\alpha} = \int_0^1 \alpha(y)s(y)dy \right)$$


$$Qs(y)(1-\alpha(y)) - A + C\bar{T} = BT^*(y) + CT^*(y) = (B+C)T^*(y)$$

$$T^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y)) - A + C\bar{T})$$

Equilibrium temperature profile:

$$T^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y)) - A + C\bar{T})$$

where $\bar{T} = \frac{1}{B} (Q(1-\bar{\alpha}) - A)$ and $\bar{\alpha} = \int_0^1 \alpha(y)s(y)dy$



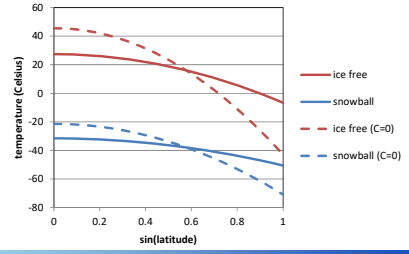
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Budyko's Model

Budyko's Equilibrium

$$Qs(y)(1-\alpha(y)) - (A + BT^*(y)) + C(\bar{T} - T^*(y)) = 0$$

Equilibrium temperature profile: $T^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y)) - A + C\bar{T})$

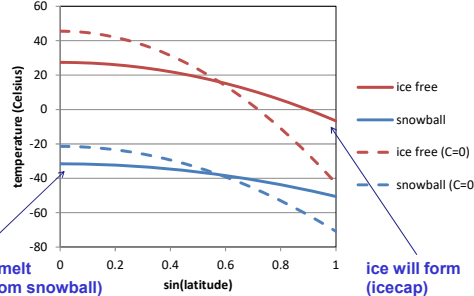


$C = 3.04$
 $\alpha(y) = 0.32$: ice free
 $\alpha(y) = 0.62$: snowball (constant albedo)

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Budyko's Model

Budyko's Equilibrium



ice won't melt (no exit from snowball) ice will form (icecap)

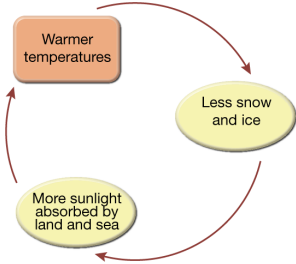
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Budyko's Model

Ice Albedo Feedback

temperature warms
ice melts
albedo decreases
more sunlight absorbed
temperature warms
REPEAT

Why would it stop?



<http://www.i-fink.com/melting-polar-ice/>

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Budyko's Model

Ice Albedo Feedback

M. I. Budyko, "The effect of solar radiation variations on the climate of the Earth," *Tellus XXI*, 611-619, 1969.

temperature warms
ice melts
albedo decreases
more sunlight absorbed
temperature warms
REPEAT

Why would it stop?



http://www.inenco.org/index_principals.html

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Budyko's Model


Ice Albedo Feedback

Why would it stop?


Budyko's Equation

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y)) - (A + BT) + C(\bar{T} - T)$$

Labels for the equation: R (heat capacity), $\frac{\partial T}{\partial t}$ (surface temperature), $Qs(y)$ (insolation), $1-\alpha(y)$ (albedo), $A + BT$ (OLR), $C(\bar{T} - T)$ (heat transport), $\bar{T} = \int_0^1 T(y)dy$ (sin(latitude)).



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Budyko's Model

Ice Albedo Feedback

What if the albedo is not constant?

Ice Line Assumption: There is a single ice line at $y=\eta$ between the equator and the pole. The albedo is α_1 below the ice line and α_2 above it.

$$\alpha(y, \eta) = \begin{cases} \alpha_1 & y < \eta \\ \alpha_2 & y > \eta \end{cases}$$

Equilibrium condition:


$$Qs(y)(1 - \alpha(y, \eta)) - (A + BT_s^*(y)) + C(\bar{T} - T_s^*(y)) = 0$$

Equilibrium solution:

$$T_s^*(y) = \frac{1}{B+C} (Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*)$$

where $\bar{T}^* = \frac{1}{B} (Q(1 - \bar{\alpha}(\eta)) - A)$ ($\bar{\alpha}(\eta) = \int_0^1 \alpha(y, \eta) s(y) dy$)

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Budyko's Model

Ice Albedo Feedback

equilibrium temperature profile:

$$T_s^*(y) = \frac{1}{B+C} (Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*), \text{ where } \bar{T}^* = \frac{1}{B} (Q(1 - \bar{\alpha}(\eta)) - A)$$

albedo:

$$\alpha(y, \eta) = \begin{cases} \alpha_1 = 0.32, & y < \eta \\ \alpha_2 = 0.62, & y > \eta \end{cases}$$

global albedo:

$$\bar{\alpha}(\eta) = \int_0^\eta \alpha_1 s(y) dy + \int_\eta^1 \alpha_2 s(y) dy$$

let:

$$S(\eta) = \int_0^\eta s(y) dy, \quad 1 - S(\eta) = \int_\eta^1 s(y) dy, \quad \text{since } 1 = \int_0^1 s(y) dy$$


then:

$$\bar{\alpha}(\eta) = \alpha_1 S(\eta) + \alpha_2 (1 - S(\eta)) = \alpha_2 - (\alpha_2 - \alpha_1) S(\eta) = 0.62 - 0.3S(\eta)$$

$$S(\eta) = \int_0^\eta s(y) dy = \int_0^\eta (1 - 0.241(3y^2 - 1)) dy = \eta - 0.241(\eta^3 - \eta)$$

Chylek & Coakley

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Budyko's Model

Ice Albedo Feedback

equilibrium temperature profile:

$$T_s^*(y) = \frac{1}{B+C} (Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*), \text{ where } \bar{T}^* = \frac{1}{B} (Q(1 - \bar{\alpha}(\eta)) - A)$$

albedo:

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$$S(\eta) = \int_0^\eta s(y) dy = \int_0^\eta (1 - 0.241(3y^2 - 1)) dy = \eta - 0.241(\eta^3 - \eta)$$

Chylek & Coakley

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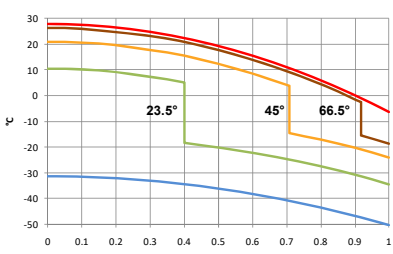


Budyko's Model


Ice Albedo Feedback

$$T_s^*(y) = \frac{1}{B+C} (Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}^*)$$

For each fixed η , there is an equilibrium solution for Budyko's equation



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Budyko's Model

Dynamics

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y)) - (A + BT) + C(\bar{T} - T)$$

Let X be the space of functions where T lives. (e.g. $L^1([0,1])$)
Let

$$L: X \rightarrow X: LT = C\bar{T} - (B+C)T,$$

$$f(y) = Qs(y)(1 - \alpha(y)) - A$$


Budyko's equation can be written as a linear vector field on X .

$$R \frac{dT}{dt} = f + LT$$

The operator L has only point spectrum, with all eigenvalues negative.
Therefore, all solutions are stable.
True for any albedo function.

experts only

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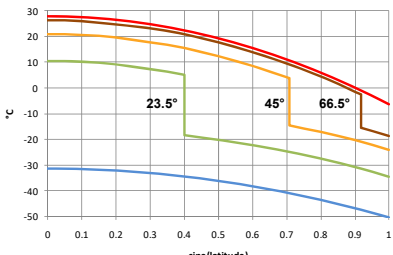
Budyko's Model

Ice Albedo Feedback


$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

For each fixed η , there is a **globally stable** equilibrium solution for Budyko's equation.

How to pick one?



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Budyko's Model


Ice Albedo Feedback

Summary


If we artificially hold the ice line at a fixed latitude, then the surface temperature will come to an equilibrium.

However, if the temperature is high, we would expect ice to melt and the ice line to retreat to higher latitudes. If the temperature is low, we would expect ice to form and the ice line to advance to lower latitudes.

How to model this expectation?



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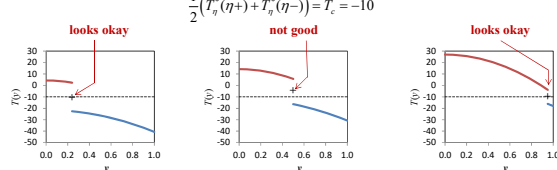

Budyko's Model

Ice Albedo Feedback


For each fixed η , there is a stable equilibrium solution for Budyko's equation.

Standard assumption: Permanent ice forms if the annual average temperature is below $T_c = -10^\circ\text{C}$ and melts if the annual average temperature is above T_c .

Additional condition: The average temperature across the ice boundary is the critical temperature T_c .

$$\frac{1}{2}(T_\eta^*(\eta^+) + T_\eta^*(\eta^-)) = T_c = -10$$



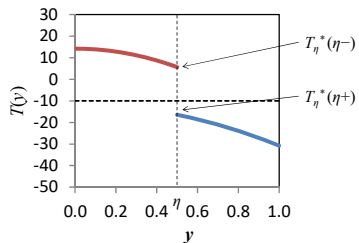

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Budyko's Model

Ice Albedo Feedback

ice line condition: $\frac{1}{2}(T_\eta^*(\eta^+) + T_\eta^*(\eta^-)) = T_c = -10$

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Budyko's Model

Ice Albedo Feedback

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

Equilibrium: $T_\eta^*(y) = \frac{1}{B+C}(Qs(y)(1 - \alpha(y, \eta)) - A + C\bar{T}_\eta^*)$


Ice line condition: $\frac{1}{2}(T_\eta^*(\eta^+) + T_\eta^*(\eta^-)) = T_c = -10$

Albedo: $\alpha(y, \eta) = \begin{cases} \alpha_1 = 0.32, & y < \eta \\ \alpha_2 = 0.62, & y > \eta \end{cases} \quad \alpha(\eta^-, \eta) = \alpha_1, \quad \alpha(\eta^+, \eta) = \alpha_2$


$$T_\eta^*(\eta^+) = \frac{1}{B+C}(Qs(\eta)(1 - \alpha_2) - A + C\bar{T}_\eta^*) \quad T_\eta^*(\eta^-) = \frac{1}{B+C}(Qs(\eta)(1 - \alpha_1) - A + C\bar{T}_\eta^*)$$

Ice line condition: $\frac{1}{2}(T_\eta^*(\eta^+) + T_\eta^*(\eta^-)) = \frac{1}{B+C}(Qs(\eta)(1 - \alpha_0) - A + C\bar{T}_\eta^*) = T_c = -10$

where: $\alpha_0 = \frac{\alpha_1 + \alpha_2}{2} = 0.47$



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Budyko's Model

Ice Albedo Feedback

$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$


Ice line condition: $\frac{1}{B+C}(Qs(\eta)(1 - \alpha_0) - A + C\bar{T}_\eta^*) = T_c = -10$

Rewrite: $h(\eta) = \frac{1}{B+C}(Qs(\eta)(1 - \alpha_0) - A + C\bar{T}_\eta^*) - T_c = 0$


Recall equilibrium GMT: $\bar{T}_\eta^* = \frac{1}{B}(Q(1 - \bar{\alpha}(\eta)) - A)$

Recall average albedo: $\bar{\alpha}(\eta) = \int_0^\eta \alpha(y, \eta) s(y) dy = \alpha_2 - (\alpha_2 - \alpha_1)S(\eta) = 0.62 - 0.3S(\eta)$

where: $S(\eta) = \int_0^\eta s(y) dy = \eta - 0.241(\eta^3 - \eta)$

$$h(\eta) = \frac{Q}{B+C} \left(s(\eta)(1 - \alpha_0) + \frac{C}{B}(1 - \alpha_2 + (\alpha_2 - \alpha_1)S(\eta)) \right) - \frac{A}{B} - T_c = 0$$


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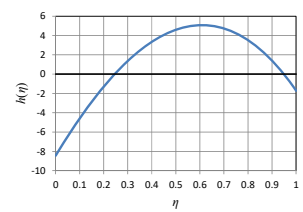
Budyko's Model

Ice Albedo Feedback


$$R \frac{\partial T}{\partial t} = Qs(y)(1 - \alpha(y, \eta)) - (A + BT) + C(\bar{T} - T)$$

The additional condition: $\frac{1}{2}(T_\eta^*(\eta^+) + T_\eta^*(\eta^-)) = T_c = -10$

can be written: $h(\eta) = \frac{Q}{B+C} \left(s(\eta)(1 - \alpha_0) + \frac{C}{B}(1 - \alpha_2 + (\alpha_2 - \alpha_1)S(\eta)) \right) - \frac{A}{B} - T_c = 0$



Two equilibria (zeros of h) satisfy the additional condition.



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Budyko's Model

Ice Albedo Feedback

Equilibrium temperature profiles

$$T_{\eta}^*(y) = \frac{1}{B+C} (Qs(y)(1-\alpha(y,\eta)) - A + C\bar{T}_{\eta}^*)$$

Interesting Solutions:

- small cap
- large cap
- ice free
- snowball

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Budyko's Model

Dynamics of the Ice Line

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y,\eta)) - (A+BT) + C(\bar{T}-T)$$

Idea:

If the average temperature across the ice line is above the critical temperature, some ice will melt, moving the ice line toward the pole. If it is below the critical temperature, the ice will advance toward the equator.

Widiasih's equation: $\frac{d\eta}{dt} = \varepsilon(T(\eta) - T_c)$

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Budyko's Model

Dynamics of the Ice Line

Budyko-Widiasih Model

$$\frac{d\eta}{dt} = \varepsilon(T(\eta) - T_c)$$

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y,\eta)) - (A+BT) + C(\bar{T}-T)$$

State space: $[0, 1] \times X$

Esther R. Widiasih, Dynamics of the Budyko Energy Balance Model, *SIAM J. Appl. Dyn. Syst.*, 12 (2013), 2068–2092.

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Budyko's Model

Dynamics of the Ice Line

$$\frac{d\eta}{dt} = \varepsilon(T(\eta) - T_c)$$

State space: $[0, 1] \times X$

$$R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y,\eta)) - (A+BT) + C(\bar{T}-T)$$

Widiasih's Theorem. For sufficiently small ε , the system has an attracting invariant curve given by the graph of a function $\Phi_{\varepsilon} : [0, 1] \rightarrow X$. On this curve, the dynamics are approximated by the equation

$$\frac{d\eta}{dt} = \varepsilon h(\eta)$$

experts only

Esther R. Widiasih, Dynamics of the Budyko Energy Balance Model, *SIAM J. Appl. Dyn. Syst.*, 12 (2013), 2068–2092.

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Budyko's Model

Budyko-Widiasih Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta)$$

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Budyko's Model

Summary

surface temperature $\bar{T} = \int_0^1 T(y) dy$

heat capacity $R \frac{\partial T}{\partial t} = Qs(y)(1-\alpha(y)) - (A+BT) + C(\bar{T}-T)$

insolation $Qs(y)$

albedo $\alpha(y)$

OLR $(A+BT)$

heat transport $C(\bar{T}-T)$

reduces to

$$\frac{d\eta}{dt} = \varepsilon h(\eta)$$

$$\frac{d\eta}{dt} = \varepsilon h(\eta) = \varepsilon \left(\frac{Q}{B+C} (s(\eta)(1-\alpha_0) + \frac{C}{B} (1-\alpha_2 + (\alpha_2 - \alpha_1) S(\eta))) - \frac{A}{B} - T_c \right)$$

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Budyko's Model

Budyko-Widiasih Model

$$R \frac{\partial T}{\partial t} = Q_s(y)(1 - \alpha(y)) - (A + BT) + C(\bar{T} - T)$$

$\bar{T} = \int_0^1 T(y) dy$

What about the greenhouse effect?
 $A + BT$ is the outgoing long wave radiation. This term decreases if the greenhouse gases increase.
 We view A as a parameter.

$$\frac{d\eta}{dt} = \epsilon h(\eta, A)$$

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Budyko's Model

Budyko-Widiasih Model

$$\frac{d\eta}{dt} = \epsilon h(\eta, A)$$

isocline $h(\eta, A) = 0$

high CO₂ low CO₂

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Snowball Earth

Is it possible for Earth to become completely covered in ice?
(Snowball Earth)

Did it ever happen?

<http://www.astrobio.net/topic/solar-system/earth/new-information-about-snowball-earth-period/>

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Snowball Earth

There is evidence that Snowball Earth has occurred, the last time about 600 million years ago.

regional glaciation global glaciation oxygen level

4.0 3.0 2.0 1.0 0 100%
 21%
 10%
 1%
 0.1%
 log pO₂

4.567 4.0 3.0 2.0 1.0 0
 Time (billions of years before present)

Solar System Origin
 Hadean Archean Proterozoic Phan.
 Late Heavy Bombardment
 bacteria & archaea algae animals, fungus, plants
 O₂ levels uncertain

<http://www.snowball-earth.org/when.html>

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Snowball Earth

The continents were clustered near the equator.

Hoffman & Schrag, *Snowball Earth*, SCIENTIFIC AMERICAN, January 2000, 68-75

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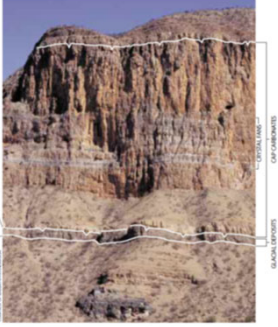
Snowball Earth

"Ice-rafted debris" occurred in ocean sediments near the equator, indicating large equatorial glaciers calving icebergs.

Hoffman & Schrag, *Snowball Earth*, SCIENTIFIC AMERICAN, January 2000, 68-75

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Snowball Earth



Large limestone deposits "cap carbonates" are found immediately above the glacial debris, indicating a rapid warming period following the snowball.

Hoffman & Schrag, *Snowball Earth*, SCIENTIFIC AMERICAN, January 2000, 68-75

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Snowball Earth

Idea:

When the Earth is mostly ice-covered, silicate weathering slows down, but volcanic activity stays the same, allowing for a build-up of CO₂ in the atmosphere.

When the Earth is mostly ice-free, silicate weathering speeds up, drawing down the CO₂ in the atmosphere.

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Budyko's Model

Budyko-Widiasih Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A)$$

What if A is a dynamical variable?

Simple equation:

$$\frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c)$$

$0 < \tilde{\varepsilon} \ll \varepsilon \ll 1$ ← MCRN Paleocarbon equation (silicate weathering)

New system:

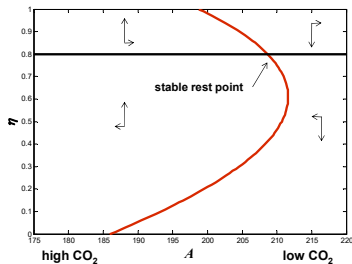
$$\begin{cases} \frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c) \\ \frac{d\eta}{dt} = \varepsilon h(\eta, A) \end{cases}$$

Anna M. Barry, Esther Widiasih, Richard McGehee, *Discrete & Continuous Dynamical Systems-B*, 22 (2017) doi: 10.3934/dcdsb.2017125

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Budyko's Model

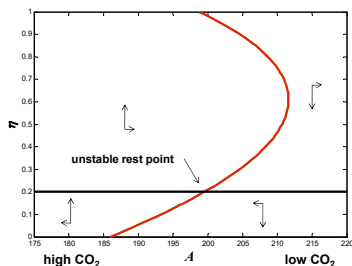
Budyko-Widiasih-Paleocarbon Model

$$\frac{d\eta}{dt} = \varepsilon h(\eta, A), \quad \frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c)$$


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Budyko's Model

Budyko-Widiasih-Paleocarbon Model

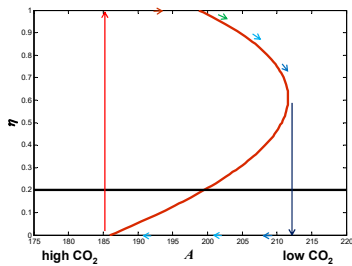
$$\frac{d\eta}{dt} = \varepsilon h(\eta, A), \quad \frac{dA}{dt} = \tilde{\varepsilon}(\eta - \eta_c)$$


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Budyko's Model

Budyko-Widiasih-Paleocarbon Model


Snowball – Hothouse Oscillations



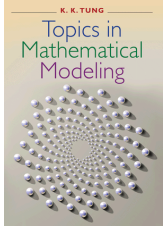
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Budyko's Model
Suggested Reading



Hoffman & Schrag, *Snowball Earth*,
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