

Moving Frames for Pseudo-Groups.

I. The Maurer–Cartan Forms

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Sur la théorie, si importante sans doute, mais pour nous si obscure, des «groupes de Lie infinis», nous ne savons rien que ce qui trouve dans les mémoires de Cartan, première exploration à travers une jungle presque impénétrable; mais celle-ci menace de refermer sur les sentiers déjà tracés, si l'on ne procède bientôt à un indispensable travail de défrichage.

— André Weil, [76]

Abstract.

This paper begins a series devoted to developing general and practical theory of moving frames for infinite-dimensional Lie pseudo-groups. In this first, preparatory part, we present a new, direct approach to the construction of invariant Maurer–Cartan forms and the Cartan structure equations for a pseudo-group. Our approach is completely explicit and avoids reliance on the theory of exterior differential systems and prolongation.

The second paper [60] will apply these constructions in order to develop the moving frame algorithm for the action of the pseudo-group on submanifolds. The third paper [61] will apply Gröbner basis methods to prove a fundamental theorem on the freeness of pseudo-group actions on jet bundles, and a constructive version of the finiteness theorem of Tresse and Kumpera for generating systems of differential invariants and also their syzygies. Applications of the moving frame method include practical algorithms for constructing complete systems of differential invariants and invariant differential forms, classifying their syzygies and recurrence relations, analyzing invariant variational principles, and solving equivalence and symmetry problems arising in geometry and physics.

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1. Introduction.

Lie pseudo-groups are the infinite-dimensional counterparts of a local Lie groups of transformations. In Lie's day, abstract Lie groups were as yet unknown, and, as a result, no significant distinction was drawn between finite-dimensional and infinite-dimensional theory. However, since then the two subjects have traveled along radically different paths. The finite-dimensional theory has been rigorously formalized, and is a well accepted and widely used mathematical tool. In contrast, the theory of infinite-dimensional pseudo-groups remains surprisingly primitive in its current overall state of development. Since there is still no generally accepted abstract object to play the role of the infinite-dimensional Lie group, Lie pseudo-groups only arise through their concrete action on a space. This makes the classification and analytical foundations of the subject problematic, particularly in the intransitive situation. We refer the reader to the original papers of Lie, Medolaghi, and Vessiot, [43, 44, 51, 74], for the classical theory of pseudo-groups, to Cartan, [13, 14], for their reformulation in terms of exterior differential systems, and [21, 22, 33, 34, 35, 41, 42, 46, 62, 63, 67, 68, 70, 72] for a variety of modern approaches. Lie pseudo-groups appear in many fundamental physical and geometrical contexts, including gauge symmetries, [6], Hamiltonian mechanics and symplectic and Poisson geometry, [54], conformal geometry of surfaces and conformal field theory, [20, 23], the geometry of real hypersurfaces, [17], symmetry groups of both linear and nonlinear partial differential equations, such as the Navier-Stokes and Kadomtsev–Petviashvili (KP) equations appearing in fluid and plasma mechanics, [5, 11, 19, 54], and geometric numerical integration, [50]. Pseudogroups also appear as foliation-preserving groups of transformations, with the associated characteristic classes defined by certain invariant forms, cf. [26].

The main goal of this series of papers is to develop a moving frame theory for Lie pseudo-groups. The motivation behind moving frames or “repères mobiles”, as developed by Cartan, [12], is to solve the equivalence problem for submanifolds under the action of a transformation group; see also [28, 27, 32]. In a series of recent collaborative papers, starting with [24, 25], the first author has successfully reformulated the classical theory in a completely general, algorithmic, equivariant framework that can be readily applied to arbitrary finite-dimensional Lie group actions. Applications have included complete classifications of differential invariants and their syzygies, equivalence and symmetry properties of submanifolds, rigidity theorems, invariant signatures in computer vision, [3, 7, 9, 57], joint invariants and joint differential invariants, [8, 57], invariant numerical algorithms, [57, 58], classical invariant theory, [4, 56], Poisson geometry and solitons, [49], and the calculus of variations, [37, 38]. New applications of these methods to computation of symmetry groups and classification of partial differential equations can be found in [48, 52].

In [24], the preliminary outlines of a comparable theory for infinite-dimensional pseudo-groups was illustrated by a few selected examples. In this paper, we shall realize the goal of a complete, rigorous moving frame theory for a wide class of pseudo-group actions. The moving frame theory will allow us to systematically extend all of the cited applications to Lie pseudo-groups, as well as providing new methods and new insights into the foundations of the subject.

Throughout, M will be a smooth manifold, and \mathcal{G} a Lie pseudo-group acting thereon. The action of \mathcal{G} on the submanifolds $S \subset M$ becomes a tale of two (infinite) jet bundles.

On the one hand, the space of jets of pseudo-group transformations forms a subbundle $\mathcal{G}^{(\infty)} \subset \mathcal{D}^{(\infty)}$ of the bundle $\mathcal{D}^{(\infty)} = \mathcal{D}^{(\infty)}(M) \subset J^\infty(M, M)$ of infinite jets (Taylor series) of local diffeomorphisms $\varphi: M \rightarrow M$. On the other hand, the action of the pseudo-group diffeomorphisms on p -dimensional submanifolds $S \subset M$ induces the prolonged pseudo-group action on the corresponding submanifold jet bundle $J^\infty(M, p)$. The ultimate goal of a classical moving frame theory is to analyze the geometry of the prolonged pseudo-group action on $J^\infty(M, p)$.

Our constructions will be based on an amalgamation of two powerful, general theories: groupoids and the variational bicomplex. As first pointed out by Ehresmann, [21], the diffeomorphism jet bundle $\mathcal{D}^{(\infty)}$ carries the structure of a groupoid, [22, 47], and $\mathcal{G}^{(\infty)}$ should be viewed as a sub-groupoid thereof. The groupoid category encodes the fact that diffeomorphism jets can only be composed when the target of the first matches the source of the second. Thus, we shall be required to adapt the Lie group based moving frame constructions arising in the finite-dimensional theory into a groupoid version that applies to more general pseudo-groups.

Following [37, 38], we shall develop the theory in the context of the variational bicomplex over each (infinite) jet bundle, [1, 2, 38, 73, 75]. The bicomplex construction[†] relies on a choice of local bundle structure on M . The interactions between the two bicomplexes associated with $J^\infty(M, p)$ and $J^\infty(M, M)$ provide the key to understanding the moving frame constructions. In particular, the invariant contact forms on $\mathcal{D}^{(\infty)} \subset J^\infty(M, M)$ will play the role of Maurer–Cartan forms for the diffeomorphism pseudo-group. This identification enables us to formulate explicitly the complete system of structure equations in power series form. Restricting the diffeomorphism-invariant forms to the pseudo-group subbundle $\mathcal{G}^{(\infty)} \subset \mathcal{D}^{(\infty)}$ yields a complete system of Maurer–Cartan forms for the pseudo-group. The remarkable fact is that the Maurer–Cartan forms satisfy an “invariantized” version of the linear infinitesimal determining equations for the pseudo-group, and, as a result, we can immediately produce an explicit form of the pseudo-group structure equations. Our direct algorithm effectively bypasses the more intricate recursive prolongation procedure of Cartan, [13, 14, 34, 45, 46]. Application of these results to the design of a practical algorithm for directly determining the structure of symmetry (pseudo-)groups of partial differential equations can be found in [15].

The moving frame theory for pseudo-groups acting on submanifold jet bundles will be developed in second paper in this series, [60], which will appear elsewhere. As in the finite-dimensional theory, [25], a moving frame for a pseudo-group action relies on the choice of a cross-section to the pseudo-group orbits. Once the moving frame is determined via normalization of the pseudo-group parameters, the resulting invariantization process leads to a complete system of differential invariants, invariant differential forms, and so on. The all-important recurrence formulae, leading to the determination of a complete,

[†] An intrinsic formulation requires the more abstract language of spectral sequences, based on the contact filtration of the cotangent bundle, [75]. However, since we shall always deal with local coordinate formulae, we have chosen to keep our constructions within the more down-to-earth bicomplex version. We refer the reader to Itskov’s thesis, [29, 30], for the more abstract \mathcal{C} -spectral sequence approach.

finite system of fundamental differential invariants and their syzygies or differential relations, will rely on the Maurer–Cartan form constructions developed here. The fact that we are eventually dealing with two distinct jet bundles and their associated bicomplexes complicates the formalism, and establishing a consistent and self-explanatory notation and terminology is half the battle. This is particularly important once we need to consider differential forms, since there are two different systems of contact forms, horizontal forms, total differentiations, etc.

More substantial applications of these results and algorithms are now beginning to appear. In [15, 16], our methods are used to analyze symmetry groups of integrable partial differential equations, specifically the Korteweg–deVries (KdV) and Kadomtsev–Petviashvili (KP) equations, while [5] develops applications to the symmetry groups of partial differential equations arising in fluid mechanics and climatology, modeling shallow water and semi-geostrophic flows. Complete theoretical justifications of our methods will be deferred until the third paper, [61], which relies on new methods from the modern theory of Gröbner bases, [18]. In particular, a new, constructive version of the Tresse–Kumpera finiteness theorem for differential invariants will be established therein.

2. Pseudo–Groups.

Let M be a smooth m -dimensional manifold, and let $\mathcal{D} = \mathcal{D}(M)$ be the pseudo-group of all local[†] diffeomorphisms $\varphi: M \rightarrow M$. For each $n \geq 0$, let $\mathcal{D}^{(n)} = \mathcal{D}^{(n)}(M) \subset \mathbf{J}^n(M, M)$ denote the bundle of their n^{th} order jets. In particular $\mathcal{D}^{(0)} = M \times M$. The natural projections are written $\pi_k^n: \mathcal{D}^{(n)} \rightarrow \mathcal{D}^{(k)}$, and we let $\mathcal{D}^{(\infty)}$ be the inverse limit — the bundle of infinite jets or Taylor series of local diffeomorphisms. Each bundle $\mathcal{D}^{(n)}$, $0 \leq n \leq \infty$, carries the structure of a groupoid, [21, 22, 47]. The *source map* $\sigma^{(n)}(\mathbf{j}_n \varphi|_z) = z$ and *target map* $\tau^{(n)}(\mathbf{j}_n \varphi|_z) = \varphi(z) = Z$ induce the double fibration

$$\begin{array}{ccc} & \mathcal{D}^{(n)} & \\ \sigma^{(n)} \swarrow & & \searrow \tau^{(n)} \\ M & & M \end{array} \quad (2.1)$$

Following Cartan, [13, 14], we will consistently use lower case letters, z, x, u, \dots for the source coordinates and the corresponding upper case letter Z, X, U, \dots for the target coordinates of our diffeomorphisms $Z = \varphi(z)$.

The groupoid multiplication follows from the composition of local diffeomorphisms. Thus, given $g^{(n)} = \mathbf{j}_n \varphi|_z$, $h^{(n)} = \mathbf{j}_n \psi|_Z$, with $Z = \tau^{(n)}(\mathbf{j}_n \varphi|_z) = \sigma^{(n)}(\mathbf{j}_n \psi|_Z)$, we set $h^{(n)} \cdot g^{(n)} = \mathbf{j}_n(\psi \circ \varphi)|_z$. Let $\mathbf{1}^{(n)}$ denote the identity section of $\sigma^{(n)}: \mathcal{D}^{(n)} \rightarrow M$, whose value at a point $z \in M$ is the n -jet $\mathbf{1}_z^{(n)}$ of the identity diffeomorphism. Local diffeomorphisms $\psi \in \mathcal{D}$ can act on $\mathcal{D}^{(n)}$ by either left or right multiplication:

$$\mathbf{L}_\psi(\mathbf{j}_n \varphi|_z) = \mathbf{j}_n(\psi \circ \varphi)|_z, \quad \mathbf{R}_\psi(\mathbf{j}_n \varphi|_z) = \mathbf{j}_n(\varphi \circ \psi^{-1})|_{\psi(z)}. \quad (2.2)$$

[†] Our notational conventions allow the domain of definition of a map $\varphi: M \rightarrow N$ to be a proper open subset: $\text{dom } \varphi \subset M$. Also, when we write $\varphi(z)$ we implicitly assume $z \in \text{dom } \varphi$.

Given $z \in M$, let $\mathcal{D}^{(n)}|_z = (\boldsymbol{\sigma}^{(n)})^{-1}\{z\}$ denote the source jet fiber. The *isotropy sub-group* $\mathcal{D}_z = \{\varphi \in \mathcal{D} \mid \varphi(z) = z\}$ consists of all local diffeomorphisms that fix z . The n^{th} order frame bundle of M , [36], $\mathcal{F}^{(n)} = \{g^{(n)} \in \mathcal{D}^{(n)} \mid \boldsymbol{\sigma}^{(n)}(g^{(n)}) = \boldsymbol{\tau}^{(n)}(g^{(n)})\}$ can be identified as the subbundle whose fibers $\mathcal{F}^{(n)}|_z = \mathcal{D}_z^{(n)} \subset \mathcal{D}^{(n)}|_z$ consist of jets $g^{(n)} = j_n \varphi|_z$ of isotropy maps $\varphi \in \mathcal{D}_z$. For $n < \infty$, $\mathcal{F}^{(n)}$ forms a principal bundle over M whose structure group is a finite-dimensional Lie group, namely the n^{th} order *prolonged linear group*

$$\text{GL}^{(n)}(m) = \{j_n \varphi|_0 \mid \varphi \in \mathcal{D}(\mathbb{R}^m), \varphi(0) = 0\} \quad (2.3)$$

prescribed by n -jets of local diffeomorphisms $\varphi: \mathbb{R}^m \rightarrow \mathbb{R}^m$ that fix the origin, [36; p. 139], [53, 59]. The group multiplication is induced by composition of diffeomorphisms.

Given local coordinates $(z, Z) = (z^1, \dots, z^m, Z^1, \dots, Z^m)$ on an open subset of $M \times M$, the induced local *Taylor coordinates* of $g^{(n)} = j_n \varphi|_z \in \mathcal{D}^{(n)}$ are denoted $(z, Z^{(n)})$, where the components Z_j^a of $Z^{(n)}$, for $a = 1, \dots, m$, $\#J \leq n$, represent the partial derivatives $\partial^J \varphi^a / \partial z^J$ of the components $\varphi^a(z^1, \dots, z^m)$ of the map φ at the source point $z = \boldsymbol{\sigma}^{(n)}(g^{(n)})$. We adopt a symmetric multi-index notation for derivatives, so $J = (j_1, \dots, j_k)$ denotes an unordered $k = \#J$ -tuple with entries $1 \leq j_\nu \leq m = \dim M$. Note that

$$\dim \mathcal{F}^{(n)} = m \binom{m+n}{m}, \quad \dim \mathcal{D}^{(n)} = m \left[1 + \binom{m+n}{m} \right]. \quad (2.4)$$

Example 2.1. Consider the case $M = \mathbb{R}$ with coordinate x . Local coordinates of $g^{(n)} = j_n \varphi|_x \in \mathcal{D}^{(n)}(\mathbb{R})$ are denoted

$$(x, X, X_x, X_{xx}, \dots, X_n),$$

where X_k corresponds to the k^{th} derivative of the diffeomorphism $X = \varphi(x)$ at the source point x . An infinite jet $g^{(\infty)} = (x, X^{(\infty)}) \in \mathcal{D}^{(\infty)}$ corresponds to the Taylor series

$$X[[h]] = X + X_x h + \frac{1}{2} X_{xx} h^2 + \dots \quad (2.5)$$

of a diffeomorphism $X = \varphi(x+h)$ at the source point x . Multiplication of diffeomorphism jets corresponds to composition (and, in the finite order case, truncation) of power series[†],

$$(X, \mathbf{X}, \mathbf{X}_X, \mathbf{X}_{XX}, \dots) \cdot (x, X, X_x, X_{xx}, \dots) = (x, \mathbf{X}, \mathbf{X}_X X_x, \mathbf{X}_X X_{xx} + \mathbf{X}_{XX} X_x^2, \dots). \quad (2.6)$$

The higher-order terms can be expressed in terms of Bell polynomials via the general Fàa-di-Bruno formula, [59, 65]. Note that the groupoid multiplication (2.6) only makes sense when the source coordinate of the first jet matches the target coordinate of the second.

Roughly, a *Lie pseudo-group* $\mathcal{G} \subset \mathcal{D}$ is a sub-pseudo-group defined by an involutive system of partial differential equations, [21, 22, 39, 62, 63, 64]. For each $0 \leq n \leq \infty$, we let $\mathcal{G}^{(n)} \subset \mathcal{D}^{(n)}$ denote the sub-groupoid consisting of all n -jets of pseudo-group diffeomorphisms. For $n < \infty$, we can regard the subbundle $\mathcal{G}^{(n)}$ as representing a system of n^{th}

[†] The reader may recoil at this XXX rated formula, but if we fail to take a consistent and systematic stance on notation at the outset, we are doomed to complete confusion before the end.

order partial differential equations, known as the *determining equations* of order n for the pseudo-group. Given $z \in M$, let $\mathcal{G}_z = \mathcal{G} \cap \mathcal{D}_z$ denote the isotropy subgroup consisting of those pseudo-group transformations that fix the point $z \in M$, and $\mathcal{G}_z^{(n)} = \mathcal{G}^{(n)} \cap \mathcal{F}^{(n)}|_z$ the subgroup of isotropy jets at z .

There is some disagreement over technical hypotheses in the definition of a Lie pseudo-group in the literature. Our version is stated as follows:

Definition 2.2. A sub-pseudo-group $\mathcal{G} \subset \mathcal{D}$ will be called a *Lie pseudo-group* if there exists $n^* \geq 1$ such that the following assumptions are satisfied for all finite $n \geq n^*$:

- (a) $\mathcal{G}^{(n)} \subset \mathcal{D}^{(n)}$ forms a smooth, embedded subbundle,
- (b) $\pi_n^{n+1}: \mathcal{G}^{(n+1)} \rightarrow \mathcal{G}^{(n)}$ is a bundle map,
- (c) every smooth local solution $Z = \varphi(z)$ to the determining system $\mathcal{G}^{(n)}$ belongs to \mathcal{G} ,
- (d) $\mathcal{G}^{(n)} = \text{pr}^{(n-n^*)} \mathcal{G}^{(n^*)}$ is obtained by prolongation.

The conditions (a–d) codify the formal integrability and local solvability requirements placed on the determining equations for the pseudo-group. In this paper, we shall not require the more technical condition of involutivity, as based on Janet–Riquier theory, [31, 66], Spencer cohomology, [62, 63], or involutive divisions, [10].

The minimal value of n^* is called the *order* of the pseudo-group, as it specifies the order of the determining system of partial differential equations that characterizes the pseudo-group transformations. We note that hypothesis (a) implies that the isotropy jets $\mathcal{G}_z^{(n)} \subset \mathcal{F}^{(n)}|_z$ form a finite-dimensional Lie group for each $z \in M$.

Remark: The modern approach to pseudo-groups, [21, 22, 39, 41, 42, 62, 70], has, by and large, been founded upon the theory of G -structures, [71]. Pseudo-groups that arise in this manner are necessarily of *constant type*, meaning that the isotropy jet groups $\mathcal{G}_z^{(n)}$ are conjugate subgroups of $\text{GL}^{(n)}(m)$. In his definition of a continuous pseudo-group, Kuranishi, [41], makes the less restrictive requirement that $\dim \mathcal{G}_z^{(n)}$ is constant on M . However, the Cartan construction of essential invariants, [13, 72], demonstrates that such an assumption is overly restrictive in the intransitive case. A key feature of our approach is that this constant dimensionality condition is *not* required, and, in fact, appears to play no significant role in the general theory!

Remark: In his thesis, [30], V. Itskov proves that, under suitable hypotheses, any non-Lie pseudo-group has the same differential invariants as a certain larger Lie pseudo-group, known as its *completion*; see also Johnson, [33; p. 317]. Thus, we are not sacrificing any significant generality by restricting our attention to Lie pseudo-groups in our study.

In local coordinates, the order n determining equations defining the pseudo-group subbundle $\mathcal{G}^{(n)}$ take the form of a system of nonlinear partial differential equations

$$F^{(n)}(z, Z^{(n)}) = 0, \tag{2.7}$$

whose local solutions $Z = \varphi(z)$ are, for any $n \geq n^*$, the pseudo-group transformations. The prolonged system $\text{pr}^{(k)} \mathcal{G}^{(n)}$ is obtained by repeatedly applying the total differential

operators[†]

$$\mathbb{D}_{z^i} = \frac{\partial}{\partial z^i} + \sum_{a=1}^m \sum_{\#J \geq 0} Z_{J,i}^a \frac{\partial}{\partial Z_J^a}, \quad i = 1, \dots, m. \quad (2.8)$$

The integrability condition (d) of the definition requires that, for all $n \geq n^*$, the system $\text{pr}^{(k)} \mathcal{G}^{(n)}$ obtained by k -fold differentiation is the same as $\mathcal{G}^{(n+k)}$. We remark that the system of partial differential equations (2.7) must satisfy the constraints that it be a *nonlinear Lie equation*, [40], ensuring that the composition of any two solutions must, where defined, also be a solution.

To illustrate our constructions, we will primarily focus on two elementary, but nevertheless interesting examples of pseudo-groups. Our methods, of course, have far wider applicability, and these examples have merely been chosen for their inherent simplicity.

Example 2.3. Let $M = \{(x, u) \mid u \neq 0\} \subset \mathbb{R}^2$. Consider the transitive Lie pseudo-group $\mathcal{G} \subset \mathcal{D}(M)$ consisting of (local) diffeomorphisms of the form

$$X = f(x), \quad U = \frac{u}{f'(x)}, \quad (2.9)$$

where $f(x) \in \mathcal{D}(\mathbb{R})$ is an arbitrary local diffeomorphism. This example was introduced by Lie, [43; p. 353], in his classification of infinite-dimensional pseudo-groups acting on the plane. The pseudo-group jets are obtained by repeated differentiation; to second order,

$$\begin{aligned} X = f, \quad U = \frac{u}{f_x}, \quad X_x = f_x, \quad X_u = 0, \quad U_x = -\frac{u f_{xx}}{f_x^2}, \quad U_u = \frac{1}{f_x}, \quad (2.10) \\ X_{xx} = f_{xx}, \quad X_{xu} = X_{uu} = 0, \quad U_{xx} = -\frac{u f_{xxx}}{f_x^2} + \frac{2u f_{xx}^2}{f_x^3}, \quad U_{xu} = -\frac{f_{xx}}{f_x^2}, \quad U_{uu} = 0, \end{aligned}$$

where $x, u, X, U, X_x, X_u, U_x, U_u, X_{xx}, \dots$ denote local coordinates on $\mathcal{D}^{(n)}(M)$. In the transformation formulae, the derivatives f, f_x, f_{xx}, \dots represent the coefficients in the Taylor series for the function $f(x+h)$ at the source point, and should be viewed as independent parameters; in particular, when not followed by (x) , the symbol $f = f_0$ represents the 0-jet or constant term in the Taylor series. Thus, the n^{th} order jet groupoid $\mathcal{G}^{(n)} \subset \mathcal{D}^{(n)}$ forms a $(n+4)$ -dimensional subbundle parametrized by the source coordinates x, u and the Taylor parameters f, f_x, f_{xx}, \dots .

The involutive system characterizing $\mathcal{G}^{(n)}$ is obtained by implicitization, and equals the $(n-1)^{\text{st}}$ prolongation of the involutive first order system

$$X_x = \frac{u}{U}, \quad X_u = 0, \quad U_u = \frac{U}{u}, \quad (2.11)$$

[†] The notation will serve to distinguish these total derivatives from those on the submanifold jet bundle $J^n(M, p)$ to be used in part II. A more intrinsic characterization appears in (3.3) below.

obtained by recursively applying the total derivatives

$$\begin{aligned}\mathbb{D}_x &= \frac{\partial}{\partial x} + X_x \frac{\partial}{\partial X} + U_x \frac{\partial}{\partial U} + X_{xx} \frac{\partial}{\partial X_x} + X_{xu} \frac{\partial}{\partial X_u} + U_{xx} \frac{\partial}{\partial U_x} + U_{xu} \frac{\partial}{\partial U_u} + \dots, \\ \mathbb{D}_u &= \frac{\partial}{\partial u} + X_u \frac{\partial}{\partial X} + U_u \frac{\partial}{\partial U} + X_{xu} \frac{\partial}{\partial X_x} + X_{uu} \frac{\partial}{\partial X_u} + U_{xu} \frac{\partial}{\partial U_x} + U_{uu} \frac{\partial}{\partial U_u} + \dots.\end{aligned}\tag{2.12}$$

For instance, its first prolongation $\text{pr}^{(1)} \mathcal{G}^{(1)} = \mathcal{G}^{(2)}$ is the six-dimensional submanifold of $\mathcal{D}^{(2)}$ defined implicitly by appending the second order conditions

$$X_{xx} = -\frac{uU_x}{U^2}, \quad X_{xu} = X_{uu} = 0, \quad U_{xu} = \frac{U_x}{u}, \quad U_{uu} = 0.\tag{2.13}$$

Example 2.4. Our second illustrative example is the Lie pseudo-group \mathcal{G} consisting of local diffeomorphisms of the form

$$X = f(x), \quad Y = f'(x)y + g(x), \quad U = u + \frac{f''(x)y + g'(x)}{f'(x)},\tag{2.14}$$

acting on $M \simeq \mathbb{R}^3$. Here $f(x) \in \mathcal{D}(\mathbb{R})$ is an arbitrary local diffeomorphism of \mathbb{R} , while $g(x) \in C^\infty(\mathbb{R})$ is an arbitrary smooth function of a single variable x . This example generalizes one of Medolaghi's pseudo-groups, [51], and illustrates some of the more subtle features in our constructions. It also serves as a simplified version of the symmetry pseudo-group of the KP equation, [15], and captures many of the features of the latter. The prolonged transformations are determined by differentiating the defining equations (2.14) with respect to x, y, u , keeping in mind that $f_y = f_u = g_y = g_u = 0$. Consequently, the subbundle $\mathcal{G}^{(n)} \subset \mathcal{D}^{(n)}$ is parametrized by the Taylor coordinates $x, y, u, f, g, f_x, g_x, f_{xx}, g_{xx}, \dots, f_{n+2}, g_{n+1}$. However, at order n , the parameter f_{n+2} is redundant, since only the particular combination $f_{n+2}y + g_{n+1}$ appears in the formulae. Eliminating the parametric coordinates, we find that for $n \geq 1$, the jet groupoid $\mathcal{G}^{(n)}$ is the $(2n + 7)$ -dimensional subbundle of $\mathcal{D}^{(n)}$ defined by prolonging the first order system

$$X_y = X_u = 0, \quad Y_y = X_x \neq 0, \quad Y_u = 0, \quad Y_x = (U - u)X_x, \quad U_u = 1,\tag{2.15}$$

through successive application of the total derivatives $\mathbb{D}_x, \mathbb{D}_y, \mathbb{D}_u$.

3. Maurer–Cartan Forms for the Diffeomorphism Groupoid.

As emphasized in [38], the moving frame calculus for finite-dimensional Lie groups relies in an essential manner on the invariant Maurer–Cartan forms. Thus, the moving frame theory for pseudo-groups requires a corresponding collection of invariant differential forms on the jet groupoids. Our constructions are motivated by the original ideas of Cartan, [13, 14, 34, 72]. However, in contrast to Cartan's recursive algorithm based upon successive prolongations of the underlying exterior differential system, we will give a direct, explicit construction of the invariant forms on the infinite diffeomorphism jet bundle $\mathcal{D}^{(\infty)}$.

Since we can identify $\mathcal{D}^{(\infty)} \subset J^\infty(M, M)$ as (an open subset of) a jet bundle over a Cartesian product manifold, there is a natural splitting of the cotangent bundle $T^*\mathcal{D}^{(\infty)}$

into horizontal and vertical (contact) components, [1, 55]. We use[†]

$$d = d_M + d_G \quad (3.1)$$

to denote the consequent splitting of the differential on $\mathcal{D}^{(\infty)}$. The subscript on the vertical differential d_G refers to the groupoid structure of $\mathcal{D}^{(\infty)}$. The interpretation is fundamental:

The vertical differential d_G on $\mathcal{D}^{(\infty)}$ should be viewed as differentiation with respect to the pseudo-group parameters, and hence plays the role of the group differential d_G in the finite-dimensional theory, [38]. Moreover, the Maurer–Cartan forms for the diffeomorphism pseudo-group will be concretely realized as the (right-)invariant contact one-forms on $\mathcal{D}^{(\infty)}$.

In terms of local coordinates $g^{(\infty)} = (z, Z^{(\infty)})$, the horizontal subbundle of $T^*\mathcal{D}^{(\infty)}$ is spanned by the one-forms $dz^1 = d_M z^1, \dots, dz^m = d_M z^m$, while the vertical subbundle is spanned by the basic *contact forms*

$$\Upsilon_J^a = d_G Z_J^a = dZ_J^a - \sum_{i=1}^m Z_{J,i}^a dz^i, \quad a = 1, \dots, m, \quad \#J \geq 0. \quad (3.2)$$

Given a differential function $F: \mathcal{D}^{(\infty)} \rightarrow \mathbb{R}$, its horizontal and vertical differentials take the form

$$d_M F = \sum_{i=1}^m \mathbb{D}_{z^i} F dz^i, \quad d_G F = \sum_{a=1}^m \sum_{\#J \geq 0} \frac{\partial F}{\partial Z_J^a} \Upsilon_J^a, \quad (3.3)$$

where \mathbb{D}_{z^i} are the total differentiation operators (2.8). In particular,

$$d_M Z_J^a = \sum_{i=1}^m Z_{J,i}^a dz^i. \quad (3.4)$$

Note that the higher-order contact forms (3.2) are obtained by iterated Lie differentiation: $\Upsilon_J^a = \mathbb{D}_z^J \Upsilon^a$, where $\mathbb{D}_z^J = \mathbb{D}_{z^{j_1}} \cdots \mathbb{D}_{z^{j_k}}$ denotes a k^{th} order total derivative.

Definition 3.1. A differential form μ on $\mathcal{D}^{(n)}$, $0 \leq n \leq \infty$, is *right-invariant* if it satisfies $R_\psi^* \mu = \mu$, where defined, for every local diffeomorphism $\psi \in \mathcal{D}$.

Remark: Morally, right-invariance of differential forms should be with respect to the right groupoid multiplication. This certainly is the case for functions. However, since the domain of a groupoid element is not an open subset of $\mathcal{D}^{(\infty)}$, defining the pull-back groupoid action on differential forms becomes problematic.

Note that the splitting of forms on $\mathcal{D}^{(\infty)}$ into horizontal and contact components is also invariant under the action of \mathcal{D} on $\mathcal{D}^{(n)}$. Hence, if μ is any right-invariant differential form, so are $d_M \mu$ and $d_G \mu$. Now, the target coordinate functions $Z^a: \mathcal{D}^{(0)} \rightarrow \mathbb{R}$ are obviously right-invariant. Therefore, by the preceding remark, their differentials

$$dZ^a = d_M Z^a + d_G Z^a \equiv \sigma^a + \mu^a, \quad a = 1, \dots, m, \quad (3.5)$$

[†] We reserve the more traditional notation d_H and d_V for the splitting of the submanifold jet bundle $J^\infty(M, p)$ to be used in part II.

split into right-invariant horizontal and contact forms. Thus, the one-forms

$$\sigma^a = d_M Z^a = \sum_{i=1}^m Z_i^a dz^i, \quad a = 1, \dots, m, \quad (3.6)$$

form an invariant horizontal coframe, while

$$\mu^a = d_G Z^a = \Upsilon^a = dZ^a - \sum_{i=1}^m Z_i^a dz^i, \quad a = 1, \dots, m, \quad (3.7)$$

are the zeroth order invariant contact forms. The total differential operators $\mathbb{D}_{Z^1}, \dots, \mathbb{D}_{Z^m}$ dual to the horizontal forms $\sigma_1, \dots, \sigma_m$ are defined by the formula

$$d_M F = \sum_{a=1}^m \mathbb{D}_{Z^a} F \sigma^a. \quad (3.8)$$

Therefore,

$$\mathbb{D}_{Z^a} = \sum_{i=1}^m w_a^i \mathbb{D}_{z^i}, \quad \text{where} \quad (w_a^i(z, Z^{(1)})) = (Z_i^a)^{-1} \quad (3.9)$$

denotes the inverse of the $m \times m$ Jacobian matrix $\nabla Z = (Z_i^a)$.

Invariance of the forms (3.6) implies that the Lie derivative of a right-invariant differential form with respect to \mathbb{D}_{Z^a} will also be right-invariant. Therefore, the higher-order invariant contact forms are

$$\mu_J^a = \mathbb{D}_Z^J \mu^a = \mathbb{D}_Z^J \Upsilon^a, \quad \text{where} \quad \mathbb{D}_Z^J = \mathbb{D}_{Z^{j_1}} \cdots \mathbb{D}_{Z^{j_k}}, \quad \begin{array}{l} a = 1, \dots, m, \\ k = \#J \geq 0. \end{array} \quad (3.10)$$

The differential operators \mathbb{D}_{Z^j} mutually commute, and so the order of differentiation is immaterial. We shall view the right-invariant contact forms $\mu^{(\infty)} = (\dots \mu_J^a \dots)$ on $\mathcal{D}^{(\infty)}$ as the *Maurer–Cartan forms* for the diffeomorphism pseudo-group. They will play exactly the same role that the ordinary Maurer–Cartan forms played in the construction of the invariant variational bicomplex in the finite-dimensional Lie group theory, [38]. The complete collection $\sigma, \mu^{(\infty)}$ in (3.6), (3.10) forms a right-invariant coframe on $\mathcal{D}^{(\infty)}$.

Example 3.2. Consider the one-dimensional case $M = \mathbb{R}$. Using the coordinate notation of Example 2.1, the first few contact forms on $\mathcal{D}^{(\infty)}$ are

$$\Upsilon = d_G X = dX - X_x dx, \quad \Upsilon_x = \mathbb{D}_x \Upsilon = dX_x - X_{xx} dx, \quad \Upsilon_{xx} = \mathbb{D}_x^2 \Upsilon = dX_{xx} - X_{xxx} dx,$$

where

$$\mathbb{D}_x = \frac{\partial}{\partial x} + X_x \frac{\partial}{\partial X} + X_{xx} \frac{\partial}{\partial X_x} + X_{xxx} \frac{\partial}{\partial X_{xx}} + \cdots \quad (3.11)$$

The basic right-invariant horizontal form is

$$\sigma = d_M X = X_x dx, \quad \text{with dual total differentiation} \quad \mathbb{D}_X = \frac{1}{X_x} \mathbb{D}_x. \quad (3.12)$$

Our algorithm produces the basic right-invariant Maurer–Cartan forms

$$\mu = \Upsilon, \quad \mu_X = \mathbb{D}_X \mu = \frac{\Upsilon_x}{X_x}, \quad \mu_{XX} = \mathbb{D}_X^2 \mu = \frac{X_x \Upsilon_{xx} - X_{xx} \Upsilon_x}{X_x^3}, \quad \dots \quad (3.13)$$

Remark: While the right-invariant forms underlie the moving frame calculus, a similar construction applies to differential forms invariant under the left action of \mathcal{D} on $\mathcal{D}^{(\infty)}$. The source coordinates z^i are left-invariant, and hence their differentials

$$dz^i = w_a^i dZ^a - \sum_{a=1}^m w_a^i \Upsilon^a, \quad i = 1, \dots, m,$$

where w_a^i is defined in (3.9), split into left invariant one-forms. A full system of left-invariant forms is obtained by further Lie differentiation:

$$\psi^i = dz^i, \quad \lambda^i = \sum_{a=1}^m w_a^i \Upsilon^a, \quad \text{and} \quad \lambda_J^i = \mathbb{D}_z^J \lambda^i. \quad (3.14)$$

In the one-dimensional situation of Example 3.2, the first few left-invariant one-forms are

$$\begin{aligned} \psi &= dx, & \lambda &= \frac{\Upsilon}{X_x} = \frac{dX - X_x dx}{X_x}, \\ \lambda_x &= \frac{X_x \Upsilon_x - X_{xx} \Upsilon}{X_x^2}, & \lambda_{xx} &= \frac{X_x^2 \Upsilon_{xx} - 2X_x X_{xx} \Upsilon_x - (X_x X_{xxx} - 2X_{xx}^2) \Upsilon}{X_x^3}. \end{aligned} \quad (3.15)$$

4. The Diffeomorphism Structure Equations.

The next step in our program is to establish the structure equations for the right-invariant Maurer–Cartan forms on the diffeomorphism groupoid $\mathcal{D}^{(\infty)}$. Instead of following Cartan’s iterative procedure based on successive prolongation of Pfaffian systems, [13], we shall present a direct approach based on a power series analysis.

Given local coordinates $z = (z^1, \dots, z^m)$ on M , we use $Z[[h]]$ to denote the vector-valued Taylor series, depending on $h = (h^1, \dots, h^m)$, of a diffeomorphism $Z = \varphi(z + h)$ at the source point $z \in M$, with components

$$Z^a[[h]] = \sum_{\#J \geq 0} \frac{1}{J!} Z_J^a h^J, \quad a = 1, \dots, m. \quad (4.1)$$

Similarly, the Taylor series $\mathbf{Z}[[H]]$ of a diffeomorphism $\mathbf{Z} = \psi(Z + H)$ based at Z has components

$$\mathbf{Z}^a[[H]] = \sum_{\#J \geq 0} \frac{1}{J!} \mathbf{Z}_J^a H^J, \quad a = 1, \dots, m, \quad (4.2)$$

with $H = (H^1, \dots, H^m)$. The groupoid multiplication on $\mathcal{D}^{(\infty)}$ is induced by composition of power series, with $\mathbf{Z}[[Z[[h]] - Z[[0]]]]$ being the Taylor series of the composition $\psi \circ \varphi$.

Let $\Upsilon[[h]]$ denote the column vector of contact form-valued power series with individual components

$$\Upsilon^a[[h]] = \sum_{\#J \geq 0} \frac{1}{J!} \Upsilon_J^a h^J, \quad a = 1, \dots, m. \quad (4.3)$$

Similarly, $\mu[[H]]$ will denote the right-invariant contact form-valued power series with

$$\mu^a[[H]] = \sum_{\#J \geq 0} \frac{1}{J!} \mu_J^a H^J, \quad a = 1, \dots, m. \quad (4.4)$$

Equations (3.10) imply that

$$\mu[[H]] = \Upsilon[[h]] \quad \text{when} \quad H = Z[[h]] - Z[[0]], \quad (4.5)$$

which can be used to recover the Maurer–Cartan forms (3.10).

To obtain the structure equations for the diffeomorphism pseudo-group, we apply the exterior derivative d (with respect to the source variables z) to $\mu[[H]]$ while keeping h fixed. First note that, from (3.2), (3.4),

$$dH = dZ[[h]] - dZ[[0]] = d_M Z[[h]] + d_G Z[[h]] - dZ[[0]] = \nabla_h Z[[h]] dz + \Upsilon[[h]] - dZ[[0]]. \quad (4.6)$$

Here

$$\nabla_h Z[[h]] = \left(\frac{\partial Z^a}{\partial h^i} [[h]] \right)$$

denotes the $m \times m$ Jacobian matrix power series obtained by differentiating $Z[[h]]$ with respect to $h = (h^1, \dots, h^m)$, while $dz = (dz^1, \dots, dz^m)^T$ is viewed as a column vector of one-forms. Using this notation, the standard structure equations

$$d\Upsilon_J^a = - \sum_{i=1}^m \Upsilon_{J,i}^a \wedge dz^i \quad (4.7)$$

for the contact forms on $\mathcal{D}^{(\infty)}$, [55], can be written in power series form as

$$d\Upsilon[[h]] = - \nabla_h \Upsilon[[h]] \wedge dz. \quad (4.8)$$

Therefore, differentiating (4.5) and using the chain rule, we find

$$d\mu[[H]] - \nabla_H \mu[[H]] \wedge dH = d(\mu[[H]]) = d\Upsilon[[h]] = - \nabla_h \Upsilon[[h]] \wedge dz. \quad (4.9)$$

It is important to note that, in the first term $d\mu[[H]]$ in (4.9), the exterior derivative only affects the invariant forms μ , whereas in the middle term $d(\mu[[H]])$ it affects both μ and H . In view of (4.6), we obtain the structure equations in the form

$$\begin{aligned} d\mu[[H]] &= \nabla_H \mu[[H]] \wedge (\nabla_h Z[[h]] dz + \Upsilon[[h]] - dZ[[0]]) - \nabla_h \Upsilon[[h]] \wedge dz \\ &= \nabla_H \mu[[H]] \wedge (\mu[[H]] - dZ[[0]]), \end{aligned} \quad (4.10)$$

where the first and last terms canceled thanks to (4.5) and the chain rule. The constant term in the final power series is

$$\mu[[0]] - dZ[[0]] = -\sigma,$$

whose components are the negatives of the invariant horizontal forms $-\sigma^a$ defined in (3.6). Therefore, their structure equations are obtained by setting $H = 0$ in (4.10):

$$d\sigma = -d\mu[[0]] = \nabla_H \mu[[0]] \wedge \sigma. \quad (4.11)$$

Equations (4.10), (4.11) constitute the complete system of structure equations for the diffeomorphism pseudo-group.

Example 4.1. When $M = \mathbb{R}$, the Maurer–Cartan form series takes the form

$$\mu + \mu_X H + \frac{1}{2} \mu_{XX} H^2 + \cdots = \mu[[H]] = \Upsilon[[h]] = \Upsilon + \Upsilon_x h + \frac{1}{2} \Upsilon_{xx} h^2 + \cdots, \quad (4.12)$$

when

$$H = X[[h]] - X[[0]] = X_x h + \frac{1}{2} X_{xx} h^2 + \frac{1}{6} X_{xxx} h^3 + \cdots.$$

Our previous formulae (3.13) for the Maurer–Cartan forms $\mu_n = \mathbb{D}_X^n \mu$ on $\mathcal{D}^{(\infty)}(\mathbb{R})$ result from equating successive powers of h in the expansion (4.12). The diffeomorphism structure equations (4.10), (4.11) are given in the form

$$d\sigma = \mu_X \wedge \sigma, \quad d\mu[[H]] = \nabla_H \mu[[H]] \wedge (\mu[[H]] - dZ[[0]]), \quad (4.13)$$

where

$$\nabla_H \mu[[H]] = \mu_X + \mu_{XX} H + \frac{1}{2} \mu_{XXX} H^2 + \cdots$$

is merely the derivative of the Maurer–Cartan form series $\mu[[H]]$ with respect to H , while

$$\mu[[H]] - dZ[[0]] = -\sigma + \mu_X H + \frac{1}{2} \mu_{XX} H^2 + \cdots.$$

The individual components satisfy

$$\begin{aligned} d\sigma &= \mu_1 \wedge \sigma, \\ d\mu_n &= -\mu_{n+1} \wedge \sigma + \sum_{i=0}^{n-1} \binom{n}{i} \mu_{i+1} \wedge \mu_{n-i} \\ &= \sigma \wedge \mu_{n+1} - \sum_{j=1}^{[(n+1)/2]} \frac{n-2j+1}{n+1} \binom{n+1}{j} \mu_j \wedge \mu_{n+1-j}. \end{aligned} \quad (4.14)$$

In this manner, we recover the general formulae established by Cartan, [13; eq. (48)], for the structure equations for the one-dimensional diffeomorphism pseudo-group $\mathcal{D}(R)$.

Example 4.2. The Maurer–Cartan forms for the planar diffeomorphism pseudo-group $\mathcal{D}(\mathbb{R}^2)$ are obtained as follows. Using the coordinates of Example 2.3, we begin with the zeroth order right-invariant forms

$$\begin{aligned} \sigma &= d_M X = X_x dx + X_u du, & \tau &= d_M U = U_x dx + U_u du, \\ \mu &= \Upsilon = d_G X = dX - X_x dx - X_u du, & \nu &= \Psi = d_G U = dU - U_x dx - U_u du. \end{aligned} \quad (4.15)$$

The higher-order right-invariant contact forms are obtained by repeatedly applying the right-invariant differential operators

$$\mathbb{D}_X = \frac{U_u \mathbb{D}_x - U_x \mathbb{D}_u}{X_x U_u - X_u U_x}, \quad \mathbb{D}_U = \frac{-X_u \mathbb{D}_x + X_x \mathbb{D}_u}{X_x U_u - X_u U_x}, \quad (4.16)$$

to μ, ν . In particular, the first order Maurer–Cartan forms are

$$\begin{aligned}\mu_X &= \frac{U_u \Upsilon_x - U_x \Upsilon_u}{X_x U_u - X_u U_x}, & \mu_U &= \frac{X_x \Upsilon_u - X_u \Upsilon_x}{X_x U_u - X_u U_x}, \\ \nu_X &= \frac{U_u \Psi_x - U_x \Psi_u}{X_x U_u - X_u U_x}, & \nu_U &= \frac{X_x \Psi_u - X_u \Psi_x}{X_x U_u - X_u U_x},\end{aligned}\tag{4.17}$$

where

$$\begin{aligned}\Upsilon_x &= dX_x - X_{xx} dx - X_{xu} du, & \Upsilon_u &= dX_u - X_{xu} dx - X_{uu} du, \\ \Psi_x &= dU_x - U_{xx} dx - U_{xu} du, & \Psi_u &= dU_u - U_{xu} dx - U_{uu} du,\end{aligned}$$

are the first order contact forms on $\mathcal{D}^{(\infty)}(\mathbb{R}^2)$. The structure equations are found via the general power series formula (4.10), which can be written in the matrix form

$$\begin{pmatrix} d\mu[[H, K]] \\ d\nu[[H, K]] \end{pmatrix} = \begin{pmatrix} \mu_H[[H, K]] & \mu_K[[H, K]] \\ \nu_H[[H, K]] & \nu_K[[H, K]] \end{pmatrix} \wedge \begin{pmatrix} \mu[[H, K]] - dX \\ \nu[[H, K]] - dU \end{pmatrix},\tag{4.18}$$

where

$$\begin{aligned}\mu[[H, K]] &= \sum_{j,k \geq 0} \frac{1}{j! k!} \mu_{j,k} H^j K^k, & \mu[[H, K]] - dX &= -\sigma + \sum_{\substack{j+k \geq 1 \\ j,k \geq 0}} \frac{1}{j! k!} \mu_{j,k} H^j K^k, \\ \nu[[H, K]] &= \sum_{j,k \geq 0} \frac{1}{j! k!} \nu_{j,k} H^j K^k, & \nu[[H, K]] - dU &= -\tau + \sum_{\substack{j+k \geq 1 \\ j,k \geq 0}} \frac{1}{j! k!} \nu_{j,k} H^j K^k,\end{aligned}$$

with $\mu_{j,k} = \mathbb{D}_X^j \mathbb{D}_U^k \mu$, $\nu_{j,k} = \mathbb{D}_X^j \mathbb{D}_U^k \nu$. In particular, the first order structure equations reduce to

$$\begin{aligned}d\mu_X &= -\mu_{XX} \wedge \sigma - \mu_{XU} \wedge \tau + \mu_U \wedge \nu_X, \\ d\sigma &= -d\mu = \mu_X \wedge \sigma + \mu_U \wedge \tau, & d\nu_X &= -\nu_{XX} \wedge \sigma - \nu_{XU} \wedge \tau + \nu_X \wedge (\mu_X - \nu_U), \\ d\tau &= -d\nu = \nu_X \wedge \sigma + \nu_U \wedge \tau, & d\mu_U &= -\mu_{XU} \wedge \sigma - \mu_{UU} \wedge \tau - \mu_U \wedge (\mu_X - \nu_U), \\ & & d\nu_U &= -\nu_{XU} \wedge \sigma - \nu_{UU} \wedge \tau + \nu_X \wedge \mu_U.\end{aligned}\tag{4.19}$$

5. Infinitesimal Generators.

Let $\mathcal{X} = \mathcal{X}(M)$ denote the space of locally defined vector fields on M , i.e., local sections of the tangent bundle TM . Let $\mathcal{X}^{(n)} = \mathbb{J}^n TM$, $0 \leq n \leq \infty$, denote the tangent n -jet bundle. Let $\mathcal{A}^{(n)}$ denote the *Lie algebroid*, [47], of the jet groupoid $\mathcal{D}^{(n)}$ which is, by definition, the space of right-invariant vector fields that are tangent to the source fibers $\mathcal{D}^{(n)}|_z$. Each Lie algebroid vector field is uniquely specified by its value on the identity section $\mathbb{1}^{(n)}$, and so we can identify $\mathcal{A}^{(n)}$ as the space of sections of the vector bundle $\mathcal{R}^{(n)} = \bigcup_{z \in M} T(\mathcal{D}^{(n)}|_z) |_{\mathbb{1}_z^{(n)}} \rightarrow M$. There is a natural Lie algebra isomorphism $\lambda^{(n)}: \mathcal{X}(M) \xrightarrow{\sim} \mathcal{A}^{(n)}$ that maps the infinitesimal generator \mathbf{v} of a local flow Φ_t acting on M to the infinitesimal generator $\mathbf{V}^{(n)} = \lambda^{(n)}(\mathbf{v})$ of the flow induced by the left action of

Φ_t on $\mathcal{D}^{(n)}$. We will call $\mathbf{V}^{(n)}$ the n^{th} order *lift* of the vector field \mathbf{v} . The infinite order case will be simply denoted by $\mathbf{V} = \boldsymbol{\lambda}(\mathbf{v})$, and the finite order cases are obtained by truncation.

In local coordinates, the lift of a vector field

$$\mathbf{v} = \sum_{a=1}^m \zeta^a(z) \frac{\partial}{\partial z^a} \quad (5.1)$$

on M is the right-invariant vector field

$$\mathbf{V} = \sum_{a=1}^m \sum_{\#J \geq 0} \mathbb{D}_z^J \zeta^a(Z) \frac{\partial}{\partial Z_J^a} \quad (5.2)$$

on $\mathcal{D}^{(\infty)}$, in which one replaces the source variables z by target variables Z in the vector field coefficients, and then differentiates with respect to z via the chain rule using the differential operators

$$\mathbb{D}_{z^i} = \sum_{a=1}^m Z_i^a \mathbb{D}_{Z^a}, \quad i = 1, \dots, m. \quad (5.3)$$

Thus, lifting essentially coincides with prolongation of a vertical vector field on $M \times M$ to the jet bundle $J^n(M, M)$, cf. [55].

Example 5.1. In the case $M = \mathbb{R}$, the lift of a vector field $\mathbf{v} = \xi(x)\partial_x$ is obtained by prolonging the associated vertical vector field $\xi(X)\partial_X$ to $\mathcal{D}^{(\infty)}$, and so

$$\begin{aligned} \mathbf{V} = & \xi \frac{\partial}{\partial X} + X_x \xi_X \frac{\partial}{\partial X_x} + (X_x^2 \xi_{XX} + X_{xx} \xi_X) \frac{\partial}{\partial X_{xx}} + \\ & + (X_x^3 \xi_{XXX} + 3X_x X_{xx} \xi_{XX} + X_{xxx} \xi_X) \frac{\partial}{\partial X_{xxx}} + \dots \end{aligned} \quad (5.4)$$

On the level of jets, we can view $\xi, \xi_X, \xi_{XX}, \dots$ as independent jet coordinates, and hence (5.4) is a linear combination of the basis Lie algebroid generators

$$\begin{aligned} & \frac{\partial}{\partial X}, \quad X_x \frac{\partial}{\partial X_x} + X_{xx} \frac{\partial}{\partial X_{xx}} + X_{xxx} \frac{\partial}{\partial X_{xxx}} + \dots, \\ & X_x^2 \frac{\partial}{\partial X_{xx}} + 3X_x X_{xx} \frac{\partial}{\partial X_{xxx}} + \dots, \quad X_x^3 \frac{\partial}{\partial X_{xxx}} + \dots, \end{aligned}$$

and so on. The general formula can be explicitly written in terms of Bell polynomials, [59, 65]. The dual basis for the Lie algebroid $\mathcal{A}^{(n)}$ relative to the basis Maurer–Cartan forms constructed above is obtained by truncating the expansions at order n .

Let $\mathcal{Z}^{(n)}$ denote the dual bundle to the vector field jet bundle $\mathcal{X}^{(n)}$, and $\mathcal{Z}^{(\infty)}$ the direct limit. (*Warning:* $\mathcal{Z}^{(n)}$ is *not* the same as the space of jets of differential one-forms on M .) We define the *lift* of a section ζ of $\mathcal{Z}^{(\infty)}$ to be the right-invariant differential form $\boldsymbol{\lambda}(\zeta)$ on $\mathcal{D}^{(\infty)}$ that vanishes on all total vector fields, and satisfies

$$\langle \boldsymbol{\lambda}(\zeta); \boldsymbol{\lambda}(\mathbf{v}) \rangle|_{g^{(\infty)}} = \langle \zeta; j_\infty \mathbf{v} \rangle|_Z \quad \text{whenever} \quad \begin{aligned} & g^{(\infty)} \in \mathcal{D}^{(\infty)}, \\ & Z = \tau^{(\infty)}(g^{(\infty)}). \end{aligned} \quad (5.5)$$

In local coordinates, each vector field jet coordinate function ζ_J^a can be viewed as a section of $\mathcal{Z}^{(\infty)}$, and is easily shown to lift to the corresponding basic right-invariant contact form

$$\mu_J^a = \lambda(\zeta_J^a). \quad (5.6)$$

More generally, any linear function of the vector field jets, $L(z, \zeta^{(n)})$, can be viewed as a section of $\mathcal{Z}^{(n)}$, whose lift

$$\lambda[L(z, \zeta^{(n)})] = L(Z, \mu^{(n)}) \quad (5.7)$$

is formally obtained by replacing the source variables z^a by their target counterparts Z^a and the vector field jet coordinates ζ_J^a by the Maurer–Cartan forms μ_J^a .

Given a pseudo-group \mathcal{G} , let $\mathfrak{g} \subset \mathcal{X}$ denote the local Lie algebra of infinitesimal generators, i.e., the set of locally defined vector fields whose flows belong to the pseudo-group. Let $\mathfrak{g}^{(n)} \subset \mathcal{X}^{(n)}$ denote their jets. In local coordinates, we can view the subbundle $\mathfrak{g}^{(n)} \subset \mathcal{X}^{(n)}$ as defining a linear system of partial differential equations

$$L^{(n)}(z, \zeta^{(n)}) = 0 \quad (5.8)$$

for the vector field coefficients, called the *linearized* or *infinitesimal determining equations* for the pseudo-group. In practice, they are constructed by linearizing the n^{th} order determining equations (2.7) at the identity:

$$L^{(n)}(z, \zeta^{(n)}) = \mathbf{V}^{(n)} [F^{(n)}(z, Z^{(n)})] \Big|_{(z, Z^{(n)}) = \mathbb{1}_z^{(n)}}. \quad (5.9)$$

The infinitesimal determining equations form a *linear Lie equation*, cf. [40], meaning that the Lie bracket of any two solutions is again a solution.

Remark: If \mathcal{G} is the symmetry group of a system of differential equations, then the linearized determining equations (5.8) are the (involutive completion of) the usual determining equations for its infinitesimal generators obtained via Lie’s algorithm, [45, 46, 54]. Infinite-dimensional pseudo-groups arise as symmetry groups of a broad range of systems, including solitons, [19], fluid mechanics, [54], oceanography and meteorology, [69], and gauge theory, [6].

By the preceding remarks, for $n \geq n^*$, the general solution to the linearized determining equations (5.8) is precisely the space of infinitesimal generators \mathfrak{g} . The system is formally integrable by virtue of condition (b) of Definition 2.2. Under reasonable hypotheses, involutivity should follow from the involutivity of the nonlinear determining equations (2.7), although a precise statement seems tricky, [41, 42]. However, local solvability of the infinitesimal determining equations does not seem to follow from our definitions so far, and so we make the following additional mild restriction.

First, the Lie algebroid $\mathcal{L}^{(n)}$ of the jet groupoid $\mathcal{G}^{(n)}$ is the set of right-invariant vector fields on $\mathcal{G}^{(n)}$ tangent to the source fibers. A Lie algebroid vector field is uniquely determined by its value at the identity section, and hence there is a one-to-one correspondence between Lie algebroid elements and sections of $\mathcal{W}^{(n)} = \bigcup_{z \in M} T(\mathcal{G}^{(n)}|_z)|_{\mathbb{1}_z^{(n)}}$, which, under the hypotheses of Definition 2.2 is a vector subbundle $\mathcal{W}^{(n)} \subset \mathcal{R}^{(n)}$ for all $n \geq n^*$.

Definition 5.2. A pseudo-group $\mathcal{G} \subset \mathcal{D}$ is *tame* if for all $z \in M$ and all $n \geq n^*$, each $\mathbf{w} \in \mathcal{W}^{(n)}|_z$ is the lift of some $\mathbf{v} \in \mathfrak{g}$, that is, $\lambda^{(n)}(\mathbf{v})|_{\mathbf{1}_z^{(n)}} = \mathbf{w}$.

Remark: All known (at least to us) examples satisfy the tameness condition, but we do not know if it is valid for all pseudo-groups satisfying the hypotheses of Definition 2.2. From now on, all pseudo-groups are assumed to be tame.

Example 5.3. The infinitesimal generators of the pseudo-group given in Example 2.3 are all vector fields of the form

$$\mathbf{v} = \xi \frac{\partial}{\partial x} + \varphi \frac{\partial}{\partial u} = a(x) \frac{\partial}{\partial x} - a'(x) u \frac{\partial}{\partial u}, \quad (5.10)$$

where $a(x)$ is an arbitrary function of x . The coefficients $\xi = a(x)$ and $\varphi = -u a'(x)$ form the general solution to the first order infinitesimal determining system

$$\xi_x = -\frac{\varphi}{u}, \quad \xi_u = 0, \quad \varphi_u = \frac{\varphi}{u}, \quad (5.11)$$

which is obtained by linearizing the determining system (2.11) at the identity, i.e., where

$$X = x, \quad U = u, \quad X_x = U_u = 1, \quad X_u = X_{xx} = X_{xu} = \dots = U_x = U_{xx} = U_{xu} = \dots = 0.$$

The system (5.11) is involutive, and hence the higher-order infinitesimal determining systems are obtained by repeatedly applying the total derivative operators

$$\begin{aligned} \mathbb{D}_x &= \frac{\partial}{\partial x} + \xi_x \frac{\partial}{\partial \xi} + \varphi_x \frac{\partial}{\partial \varphi} + \xi_{xx} \frac{\partial}{\partial \xi_x} + \xi_{xu} \frac{\partial}{\partial \xi_u} + \varphi_{xx} \frac{\partial}{\partial \varphi_x} + \varphi_{xu} \frac{\partial}{\partial \varphi_u} + \dots, \\ \mathbb{D}_u &= \frac{\partial}{\partial u} + \xi_u \frac{\partial}{\partial \xi} + \varphi_u \frac{\partial}{\partial \varphi} + \xi_{xu} \frac{\partial}{\partial \xi_x} + \xi_{uu} \frac{\partial}{\partial \xi_u} + \varphi_{xu} \frac{\partial}{\partial \varphi_x} + \varphi_{uu} \frac{\partial}{\partial \varphi_u} + \dots. \end{aligned}$$

Example 5.4. Similarly, the infinitesimal generators of the pseudo-group (2.14) are

$$\mathbf{v} = \xi \frac{\partial}{\partial x} + \eta \frac{\partial}{\partial y} + \varphi \frac{\partial}{\partial u} = a(x) \frac{\partial}{\partial x} + [a'(x)y + b(x)] \frac{\partial}{\partial y} + [a''(x)y + b'(x)] \frac{\partial}{\partial u}, \quad (5.12)$$

where $a(x), b(x)$ are arbitrary functions. These form the general solution to the first order infinitesimal determining system

$$\xi_x = \eta_y, \quad \xi_y = \xi_u = \eta_u = \varphi_u = 0, \quad \eta_x = \varphi, \quad (5.13)$$

which is obtained by linearizing the determining system (2.15) at the identity jet. Again, the linearized determining equations are involutive at order 1.

6. The Structure Equations for a Pseudo-Group.

A complete system of right-invariant one-forms on the infinite jet groupoid $\mathcal{G}^{(\infty)} \subset \mathcal{D}^{(\infty)}$ is obtained by restricting (pulling back) the Maurer–Cartan forms (3.6), (3.10). For simplicity, we continue to denote these forms by $\sigma = (\dots \sigma^a \dots)$, $\mu^{(\infty)} = (\dots \mu_J^a \dots)$. The restricted Maurer–Cartan forms are, of course, no longer linearly independent, but are subject to the constraints prescribed by the pseudo-group. Remarkably, these constraints can be explicitly characterized by an invariant version of the linearized determining equations (5.8). The proof, though, is a simple unwinding of definitions.

Theorem 6.1. *The linear system*

$$L^{(n)}(Z, \mu^{(n)}) = 0, \quad (6.1)$$

obtained by lifting the linear determining equations (5.8) as in (5.7) serves to define the complete set of linear dependencies among the right-invariant Maurer–Cartan forms $\mu^{(n)}$ on $\mathcal{G}^{(n)}$.

Proof: At each point $z \in M$, the linear determining equations are defined by a system of linear functions on the vector field jet bundle $\mathcal{X}^{(n)}$, which we can view locally as a collection of sections of the dual bundle $\mathcal{Z}^{(n)}$. Their lift, as in (5.7), forms a system of right-invariant forms on the diffeomorphism groupoid. Since the linear determining equations annihilate the tangent spaces $T(\mathcal{G}^{(n)}|_z)|_{\mathbb{1}_z^{(n)}}$, formula (5.5) immediately implies that the tangent space to $\mathcal{G}^{(n)}$ is the annihilator of the lifted system of differential forms, and hence (6.1) follows. *Q.E.D.*

Corollary 6.2. *The structure equations for the pseudo-group are obtained by restriction of the diffeomorphism structure equations (4.10) to the kernel of the linearized involutive system, as defined by (6.1).*

An invariant coframe on the infinite pseudo-group jet bundle $\mathcal{G}^{(\infty)}$ can be identified as a basis for the solution space to the infinite prolongation of the linear system (6.1). Under appropriate hypotheses, truncating at the order of involutivity of the linearized determining system (5.8) yields an involutive Pfaffian system that serves to characterize the pseudo-group.

Example 6.3. We first illustrate the method by first treating a standard finite-dimensional Lie group action, also analyzed by Cartan, [14]. The linear fractional transformations $X = (\alpha x + \beta)/(\gamma x + \delta)$ define an action of the projective group $\text{PGL}(2, \mathbb{R})$ on $M = \mathbb{RP}^1$. The determining equations are obtained by prolonging the well-known Schwarzian equation

$$X_{xxx} X_x - \frac{3}{2} X_{xx}^2 = 0. \quad (6.2)$$

The infinitesimal generators $\mathbf{v} = \xi(x) \partial_x = (a + bx + cx^2) \partial_x$ solve the linear determining equation

$$\xi_{xxx} = 0$$

obtained by linearizing (6.2) at the identity $X = x, X_x = 1, X_{xx} = X_{xxx} = \dots = 0$.

The Maurer–Cartan forms on the jet groupoid $\mathcal{G}^{(\infty)} \subset \mathcal{D}^{(\infty)}$ are obtained by restriction of the right-invariant coframe $\sigma, \mu, \mu_X, \mu_{XX}, \dots$ from Example 3.2 to the group manifold, as defined by (6.2) and its prolongations. According to Theorem 6.1, they satisfy

$$0 = \mu_{XXX} = \mu_{XXXX} = \dots, \quad (6.3)$$

which can be verified by direct computation. Therefore, referring to (3.13), the one-forms

$$\begin{aligned} \sigma &= X_x dx, & \mu &= dX - X_x dx, & \mu_X &= \frac{\Upsilon_x}{X_x} = \frac{dX_x - X_{xx} dx}{X_x}, \\ \mu_{XX} &= \frac{X_x \Upsilon_{xx} - X_{xx} \Upsilon_x}{X_x^3} = \frac{X_x dX_{xx} - X_{xx} dX_x - \frac{1}{2} X_{xx}^2 dx}{X_x^3}, \end{aligned} \quad (6.4)$$

provide a right-invariant coframe on $\mathcal{G}^{(2)} \simeq \mathcal{G}^{(\infty)}$, with local coordinates x, X, X_x, X_{xx} . The structure equations are found by truncating (4.14) and invoking (6.3):

$$d\sigma = -d\mu = \mu_X \wedge \sigma, \quad d\mu_X = \sigma \wedge \mu_{XX}, \quad d\mu_{XX} = -\mu_X \wedge \mu_{XX}, \quad (6.5)$$

and so the one-forms σ, μ_X, μ_{XX} satisfy the structure equations of $\text{PGL}(2, \mathbb{R})$. The additional one-form μ arises from the fact that $\mathcal{G}^{(2)}$ can be identified as a principal $\text{PGL}(2, \mathbb{R})$ bundle over \mathbb{RP}^1 . Indeed, by transitivity, each target fiber $(\tau^{(2)})^{-1}\{c\} = \{X = c\} \subset \mathcal{G}^{(2)}$ can be identified with $\text{PGL}(2, \mathbb{R})$. Since $\sigma + \mu = dX = 0$ on the target fibers, the restriction of the invariant one-forms yields a complete system of (fiber-dependent linear combinations of) the standard Maurer–Cartan coframe for $\text{PGL}(2, \mathbb{R})$.

Example 6.4. The Maurer–Cartan forms for the pseudo-group (2.9) are obtained by restricting the right-invariant forms (4.15), (4.17) to the subbundle defined by the pseudo-group determining equations (2.11), (2.13). The basic invariant one-forms are

$$\begin{aligned} \sigma &= \frac{u}{U} dx = f_x dx, & \tau &= U_x dx + \frac{U}{u} du = \frac{-u f_{xx} dx + f_x du}{f_x^2}, \\ \mu &= dX - \frac{u}{U} dx = \Phi, & \nu &= dU - U_x dx - \frac{U}{u} du = -\frac{u \Phi_x}{f_x^2}, \end{aligned} \quad (6.6)$$

where

$$\Phi = df - f_x dx, \quad \Phi_x = df_x - f_{xx} dx, \quad \dots \quad \Phi_n = df_n - f_{n+1} dx, \quad (6.7)$$

are contact forms in the pseudo-group parameters. The higher-order Maurer–Cartan forms are obtained by similarly restricting the higher-order diffeomorphism Maurer–Cartan forms $(\sigma, \mu^{(\infty)})$, or, alternatively, by successively differentiating μ, ν using the dual total differential operators

$$\mathbb{D}_X = \frac{U}{u} \mathbb{D}_x - U_x \mathbb{D}_u = \frac{1}{f_x} \mathbb{D}_x + \frac{u f_{xx}}{f_x^2} \mathbb{D}_u, \quad \mathbb{D}_U = \frac{u}{U} \mathbb{D}_u = f_x \mathbb{D}_u,$$

where, on the pseudo-group bundle,

$$\mathbb{D}_x = \frac{\partial}{\partial x} + \sum_{i \geq 0} f_{i+1} \frac{\partial}{\partial f_i}, \quad \mathbb{D}_u = \frac{\partial}{\partial u}.$$

Thus, the first order Maurer–Cartan forms are

$$\begin{aligned} \mu_X &= \frac{du}{u} - \frac{dU - U_x dx}{U} = \frac{\Phi_x}{f_x}, & \nu_U &= -\frac{du}{u} + \frac{dU - U_x dx}{U} = -\frac{\Phi_x}{f_x}, \\ \mu_U &= 0, & \nu_X &= \frac{U(dU_x - U_{xx} dx) - U_x(dU - U_x dx)}{u} = \frac{u f_{xx} \Phi_x - u f_x \Phi_{xx}}{f_x^4}. \end{aligned} \quad (6.8)$$

Note that, when written in terms of the parameters f, f_x, f_{xx} , the Maurer–Cartan forms are certain linear combinations

$$\mu, \quad \nu = -U \mu_X, \quad \mu_X, \quad \mu_U = 0, \quad \nu_X = -U \mu_{XX}, \quad \nu_U = -\mu_X, \quad (6.9)$$

of the Maurer–Cartan forms (3.13) for the diffeomorphism group $\mathcal{D}(\mathbb{R})$. In Cartan’s terminology, the pseudo-group is a “holohedric prolongation” of $\mathcal{D}(\mathbb{R})$, and this correspondence is maintained on the level of their Maurer–Cartan forms.

In accordance with Theorem 6.1, the Maurer–Cartan forms (6.6), (6.8) are constrained by the invariantized version

$$\mu_X = -\frac{\nu}{U}, \quad \mu_U = 0, \quad \nu_U = \frac{\nu}{U}, \quad (6.10)$$

of the linear determining equations (5.11). The higher-order constraints are obtained by (invariant) prolongation — i.e., differentiation of (6.10) with respect to X, U . Thus, at second order

$$\mu_{XX} = -\nu_{XU} = -\frac{\nu_X}{U}, \quad \mu_{XU} = \mu_{UU} = \nu_{UU} = 0. \quad (6.11)$$

The remaining independent one-forms $\sigma, \tau, \mu, \nu, \nu_X, \nu_{XX}, \nu_{XXX}, \dots$ form a right-invariant coframe on $\mathcal{G}^{(\infty)}$. The structure equations (at order 1)

$$\begin{aligned} d\sigma &= -d\mu = -\frac{\nu \wedge \sigma}{U}, & d\mu_X &= \frac{\nu_X \wedge \sigma}{U}, \\ d\tau &= -d\nu = \nu_X \wedge \sigma + \frac{\nu \wedge \tau}{U}, & d\nu_X &= -\nu_{XX} \wedge \sigma - \frac{\nu_X \wedge (\tau + 2\nu)}{U}, \end{aligned}$$

are obtained by restriction of the structure equations (4.19). Higher order structure equations can be obtained from the general power series expansion. Indeed, the solution to the invariant determining equations (6.10) can be written in terms of the $\mathcal{D}(\mathbb{R})$ Maurer–Cartan form power series (4.12):

$$\mu[[H, K]] = \mu[[H]], \quad \nu[[H, K]] = -(U + K)\mu_H[[H]]. \quad (6.12)$$

Substituting (6.12) reduces the structure equations (4.18) to the structure equations (4.14) for $\mathcal{D}(\mathbb{R})$, in accordance with the holohedric equivalence of the two pseudo-groups.

Example 6.5. The Maurer–Cartan forms for the pseudo-group (2.14) are obtained as follows. To simplify the final expressions, we use $f(x)$ and

$$e(x, y) = f'(x)y + g(x) \quad (6.13)$$

to denote the arbitrary functions, which are subject to the constraints $e_y = f_x, f_y = 0$. For all $n \geq 1$, a complete system of local coordinates on $\mathcal{G}^{(n)}$ is given by $(x, y, u, f, e, f_x, e_x, \dots, f_{n+1}, e_{n+1})$, where

$$f_k = \partial_x^k f, \quad e_k = \partial_x^k e = y \partial_x^{k+1} f + \partial_x^k g, \quad k \geq 0.$$

The basic invariant horizontal forms are

$$\begin{aligned} \sigma &= d_M X = f_x dx, & \tilde{\sigma} &= d_M Y = f_x dy + e_x dx, \\ \tau &= d_M U = du + \frac{(f_x e_{xx} - f_{xx} e_x) dx}{f_x^2} + \frac{f_{xx} dy}{f_x}. \end{aligned} \quad (6.14)$$

The order zero invariant contact forms are

$$\mu = d_G X = \Phi, \quad \tilde{\mu} = d_G Y = \Psi, \quad \nu = d_G U = \frac{f_x \Psi_x - e_x \Phi_x}{f_x^2}, \quad (6.15)$$

where Φ, Φ_x are given in (6.7), while

$$\begin{aligned} \Psi &= de - e_x dx - f_x dy, & \Psi_x &= de_x - e_{xx} dx - f_{xx} dy, & \dots \\ \Psi_n &= de_n - e_{n+1} dx - f_{n+1} dy. \end{aligned} \quad (6.16)$$

The invariant differential operators dual to (6.14) are given by

$$\mathbb{D}_X = \frac{1}{f_x} \mathbb{D}_x - \frac{e_x}{f_x^2} \mathbb{D}_y + \frac{2f_{xx}e_x - f_x e_{xx}}{f_x^3} \mathbb{D}_u, \quad \mathbb{D}_Y = \frac{1}{f_x} \mathbb{D}_y - \frac{f_{xx}}{f_x^2} \mathbb{D}_u, \quad \mathbb{D}_U = \mathbb{D}_u, \quad (6.17)$$

where, when written in terms of the pseudo-group parameters,

$$\mathbb{D}_x = \frac{\partial}{\partial x} + \sum_{n=0}^{\infty} \left(f_{n+1} \frac{\partial}{\partial f_n} + e_{n+1} \frac{\partial}{\partial e_n} \right), \quad \mathbb{D}_y = \frac{\partial}{\partial y} + \sum_{n=0}^{\infty} f_{n+1} \frac{\partial}{\partial e_n}, \quad \mathbb{D}_u = \frac{\partial}{\partial u}.$$

Applying the differential operators (6.17) repeatedly to the order zero forms (6.15) produces the higher-order Maurer–Cartan forms

$$\begin{aligned} \mu_Y = \mu_U = \tilde{\mu}_U = \nu_U = 0, & \quad \mu_X = \tilde{\mu}_Y = \frac{\Phi_x}{f_x}, & \quad \tilde{\mu}_X = \frac{f_x \Psi_x - e_x \Phi_x}{f_x^2}, \\ \nu_X = \frac{f_x^2 \Psi_{xx} - f_x f_{xx} \Psi_x - 2f_x e_x \Phi_{xx} + (3e_x f_{xx} - e_{xx} f_x) \Phi_x}{f_x^4}, & \quad \nu_Y = \frac{f_x \Phi_{xx} - f_{xx} \Phi_x}{f_x^3}, \end{aligned}$$

and so on. In accordance with Theorem 6.1, the contact Maurer–Cartan forms satisfy the invariant linearized determining equations for the pseudo-group \mathcal{G} ; to first order,

$$\mu_X = \tilde{\mu}_Y, \quad \nu = \tilde{\mu}_X, \quad \mu_Y = \mu_U = \tilde{\mu}_U = 0.$$

By repeated differentiation, setting $\mu_{j,k,l} = \mathbb{D}_X^j \mathbb{D}_Y^k \mathbb{D}_U^l \mu$, etc., we find that

$$\tilde{\mu}_{j,1,0} = \mu_{j+1,0,0}, \quad \nu_{j,0,0} = \tilde{\mu}_{j+1,0,0}, \quad \nu_{j,1,0} = \mu_{j+2,0,0}, \quad j, k, l \geq 0, \quad (6.18)$$

and all others are zero. As a result, the one-forms

$$\sigma, \quad \tilde{\sigma}, \quad \tau, \quad \mu_n = \mu_{n,0,0}, \quad \tilde{\mu}_n = \tilde{\mu}_{n,0,0} \quad \text{for} \quad n = 0, 1, 2, \dots,$$

form a \mathcal{G} -invariant coframe on $\mathcal{G}^{(\infty)}$.

The structure equations are obtained by substituting the expansions

$$\begin{aligned} \mu[[H, K]] &= \sum_{n=0}^{\infty} \frac{1}{n!} \mu_n H^n, & \tilde{\mu}[[H, K]] &= \sum_{n=0}^{\infty} \frac{1}{n!} \tilde{\mu}_n H^n + K \sum_{n=0}^{\infty} \frac{1}{n!} \mu_{n+1} H^n, \\ \nu[[H, K]] &= \sum_{n=0}^{\infty} \frac{1}{n!} \tilde{\mu}_{n+1} H^n + K \sum_{n=0}^{\infty} \frac{1}{n!} \mu_{n+2} H^n, \end{aligned}$$

into (4.10). Those involving μ_n reduce to the structure equations (4.14) for $\mathcal{D}(R)$ while

$$d\tilde{\mu}_n = \sigma \wedge \tilde{\mu}_{n+1} + \tilde{\sigma} \wedge \mu_{n+1} + \sum_{j=0}^{n-1} \left[\binom{n}{j} - \binom{n}{j+1} \right] \tilde{\mu}_{j+1} \wedge \mu_{n-j}. \quad (6.19)$$

Moreover,

$$\begin{aligned} d\sigma &= -d\mu = -\sigma \wedge \mu_X, \\ d\tilde{\sigma} &= -d\tilde{\mu} = -\sigma \wedge \tilde{\mu}_X - \tilde{\sigma} \wedge \mu_X, \\ d\tau &= -d\nu = -d\tilde{\mu}_X = -\sigma \wedge \tilde{\mu}_{XX} - \tilde{\sigma} \wedge \mu_{XX}. \end{aligned} \quad (6.20)$$

Example 6.6. The following relatively trivial, but nevertheless interesting finite-dimensional abelian transformation group

$$X = x, \quad Y = y, \quad U = u + ax + by, \quad \text{where } (a, b) \in \mathbb{R}^2, \quad (6.21)$$

appears in Cartan, [14; p. 1357], as an illustration of his theory of essential invariants; see also [72; p. 407]. The first order determining equations are

$$X = x, \quad Y = y, \quad U = u + xU_x + yU_y, \quad X_x = Y_y = U_u = 1, \quad X_y = X_u = Y_x = Y_u = 0.$$

The jet coordinates $U_x = a, U_y = b$, coincide with the group parameters, while all higher-order derivatives are 0. The invariant horizontal forms are

$$\sigma = d_M X = dx, \quad \tilde{\sigma} = d_M Y = dy, \quad \tau = d_M U = du + U_x dx + U_y dy,$$

while the non-zero invariant contact forms are

$$\nu = d_G U = x dU_x + y dU_y, \quad \nu_X = dU_x, \quad \nu_Y = dU_y.$$

They clearly satisfy the infinitesimal determining equation

$$\nu = X\nu_X + Y\nu_Y, \quad (6.22)$$

while all higher-order Maurer–Cartan forms vanish when restricted to $\mathcal{G}^{(\infty)}$. The structure equations reduce to

$$d\sigma = d\tilde{\sigma} = d\nu_X = d\nu_Y = 0, \quad d\tau = -d\nu = \nu_X \wedge \sigma + \nu_Y \wedge \tilde{\sigma}. \quad (6.23)$$

However, owing to (6.22), only 5 of the Maurer–Cartan forms are independent, and, on the submanifold where $Y = y \neq 0$, we have

$$d\tau = -d\nu = \frac{\nu \wedge \tilde{\sigma}}{Y} + \nu_X \wedge \left(\sigma - \frac{X}{Y} \tilde{\sigma} \right), \quad (6.24)$$

reproducing the structure equations found by Cartan. The function $I = X/Y = x/y$ that appears in the last structure equation (6.24) is known as an *essential invariant*. Essential invariants depend upon the fibration of the jet bundle, and are preserved under equivalence maps on the base manifold. Indeed, the alternative coframe $\sigma, \tilde{\sigma}, \tau, \nu_X, \nu_Y$ does not introduce any essential invariant in the structure equations, because the prolonged action on the five-dimensional subbundle $\mathcal{G}^{(1)} \subset \mathcal{D}^{(1)}$ is equivalent, under a non-projectable change of variables, to the trivial translation action of \mathbb{R}^2 on \mathbb{R}^5 .

Remark: Interestingly, the “enhanced” action $U = u + ax + by + c$ has no essential invariants; the structure equations are the same, but there is no linear dependence among the Maurer–Cartan forms ν, ν_X, ν_Y that gives rise to the essential invariant.

The existence of essential invariants does not affect the moving frame algorithm to be presented in [60]. The reader may wish to use Cartan’s simple example as a test case.

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