# 3.5.11e,f

(e) doesn't pass test;  $\star$  (f) uniformly convergent.

#### 3.5.21 a,c,e

- (a) The periodic extension is not continuous, and so the best one could hope for is  $a_k, b_k \to 0$  like 1/k. Indeed,  $a_0 = -2\pi$ ,  $a_k = 0$ ,  $b_k = (-1)^{k+1}2/k$ , for k > 0.
- (c) The periodic extension is  $C^0$ , and so we expect  $a_k, b_k \to 0$  like  $1/k^2$ . Indeed,  $a_0 = \frac{2}{3}\pi^2, \ a_k = (-1)^k 4/k^2, \ b_k = 0, \ \text{for } k > 0.$
- (e) The periodic extension is  $C^{\infty}$ , and so we expect  $a_k, b_k \to 0$  faster than any (negative) power of k. Indeed,  $a_0 = 1, \ a_2 = -\frac{1}{2}$ , and all other  $a_k = b_k = 0$ .

## 3.5.22 a,f

### 3.5.26 c,e

 $\star$  (c) converges in norm;  $\star$  (e) does not converge in norm.

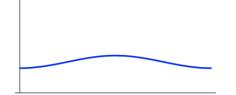
### 3.5.43

#### 4.1.7

(a) 
$$u(t,x) = \frac{1}{4} - \frac{2}{\pi^2} \sum_{j=0}^{\infty} \frac{1}{(2j+1)^2} \exp\left(-(4j+2)^2 \pi^2 t\right) \cos(4j+2)\pi x;$$
 (b)  $\frac{1}{4}$ ;

- (c) At an exponential rate of  $e^{-4\pi^2t}$ ;
- (d) As  $t \to \infty$ , the solution becomes a vanishingly small cosine wave centered around  $u = \frac{1}{4}$ , namely

$$u(t,x) \approx \frac{1}{4} - \frac{2}{\pi^2} e^{-4\pi^2 t} \cos 2\pi x$$
:



## 4.1.10 a, c

4.1.10. (a)  $u(t,x) = e^{-t} \cos x$ ; equilibrium temperature:  $u(t,x) \to 0$ .

$$\star (c) \ u(t,x) = \frac{1}{2}\pi - \frac{4}{\pi} \sum_{k=0}^{\infty} \frac{e^{-(2k+1)^2 t} \cos(2k+1)x}{(2k+1)^2}; \text{ equilibrium temperature: } u(t,x) \to \frac{1}{2}\pi.$$

# 4.1.16 a, b

(a) If  $u(t,x) = e^{\alpha t} v(t,x)$ , then

$$rac{\partial u}{\partial t} = lpha \, e^{lpha \, t} \, v(t,x) + e^{lpha \, t} \, rac{\partial v}{\partial t} \, (t,x) = \gamma \, e^{lpha \, t} \, rac{\partial^2 v}{\partial x^2} = \gamma \, rac{\partial^2 u}{\partial x^2} \, .$$

(a) If 
$$u(t,x) = e^{-\alpha t} v(t,x)$$
, then
$$\frac{\partial u}{\partial t} = \alpha e^{\alpha t} v(t,x) + e^{\alpha t} \frac{\partial v}{\partial t}(t,x) = \gamma e^{\alpha t} \frac{\partial^2 v}{\partial x^2} = \gamma \frac{\partial^2 u}{\partial x^2}.$$
(b)  $v(t,x) = e^{-\alpha t} \sum_{n=1}^{\infty} b_n e^{-(\alpha + \gamma n^2 \pi^2)t} \sin n\pi x$ , where  $b_n = 2 \int_0^1 f(x) \sin n\pi x \, dx$ 

are the Fourier sine coefficients of the initial data. All solutions tend to the equilibrium value  $u(t,x) \to 0$  as  $t \to \infty$  at an exponential rate. For most initial data, i.e., those with  $b_1 \neq 0$ , the decay rate is  $e^{-at}$ , where  $a = \alpha + \gamma \pi^2$ ; other solutions decay at a faster rate.